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# Wind Turbine Model-test Method for Achieving Similarity of Both Model and Full-Scale Thrusts and Torques

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#### 11 Abstract

12

13 For model tests of a floating offshore wind turbine (FOWT) system, a great challenge is how 14 to model the interaction between the wind turbine and floating platform with correctly-scaled 15 aerodynamic and hydrodynamic loads because of a well-known contradiction between 16 Reynolds and Froude scaling. Several approaches have been proposed in the literature to tackle 17 the challenge but none of them can correctly and simultaneously model the scaled thrust and 18 torque, and so the interaction between turbine and platform. This paper will present a new 19 model-test method for achieving similarity of both thrust and torque. This is achieved by 20 redesigning the model blades with keeping the blade twist angle same as that of the full-scale 21 turbine and by adjusting the pitch angle of blades and rotational speed of wind turbine in model 22 tests. Numerical simulations and wind tunnel model tests are carried out to validate the present 23 method. Both numerical results and experimental data show that the present method can realize the similarity of thrust and torque simultaneously, and thus, making it possible to study full 24 25 interaction between turbine and floating platform by physical testing in wind-wave basins. 26 **Keywords:** Wind turbine model test; Scaled thrust and torque; Reynolds and Froude similarity;

27 FOWT

#### 28 1. Introduction

29

The concept of FOWT (Heronemus, 1972) is a highly complex system due to the two-way coupling between hydrodynamic and aerodynamic loads on the platform and the turbine respectively, as illustrated in Fig. 1. In model tests, aerodynamic loads should be scaled based on Reynolds similarity, while hydrodynamic loads should be scaled by Froude similarity. For a FOWT model test, the contradiction between the two scaling laws makes it impossible to achieve the correctly-scaled aerodynamic and hydrodynamic loads simultaneously.

The Froude similarity is commonly used in the model tests of traditional offshore structures
(Chakrabarti, 1998; Faltinsen, 1993; Sarpkaya et al., 1981), which results in a much smaller

- **38** Reynolds number in model scale than that in full scale. Therefore, the aerodynamic loads cannot
- 39 be correctly scaled if the traditional Froude scaling method is applied to FOWT model test
- 40 directly (Martin et al., 2014; Wen et al., 2022; Madsen et al., 2020). Therefore, a key challenge
- 41 in the wind-wave basin model test of FOWT system is how to achieve the scaled aerodynamic
- 42 loads correctly with a smaller Reynolds number and thus to model the interaction between
- 43 turbine and floating platform in a correct way.



Fig. 1. Interaction between hydrodynamic and aerodynamic loads on a FOWT

There have been many studies to address the challenge and several approaches have been reviewed by Stewart et al. (2016), Chen et al. (2020) and Day et al. (2015). Some of them are summarized in Table 1.

49 In the early studies, the thrust force of wind turbine is simply simulated by a steady force on 50 the top of tower. Utsunomiya et al. (2009), by using this method, tested the motion of a sparbuoy FOWT in regular and irregular waves at NMRI (National Maritime Research Institute). 51 52 The method was also used by Zhao (2012) to study a semi-submersible FOWT. The steady 53 force method can only simulate the steady thrust and so the steady displacement of platform 54 and mooring offsets (Day et al., 2015). Clearly, this method discards the important dynamic feature of aerodynamics, and thus it only partially considers the action of the turbine on the 55 56 platform.

Roddier et al. (2010) proposed a model test method using a drag disk. In their tests on
WindFloat, the desired thrust was achieved by using a large disk on the top of tower. The
method was also employed by Wan et al. (2015; 2016; 2016; 2017) and Gao et al. (2016). This
method may be able to partially model the unsteady thrust but entirely ignores the torque.

61 Table 1

62 Approaches to simulate scaled aerodynamic loads in model tests

Approaches	Author	Similarity		Interaction between
		Thrust	Torque	turbine and platform
Steady force	Utsunomiya et al. (2009) Zhao (2012)	Partial	No	Partial effects of turbine on platform
Drag disk	Roddier et al. (2010) Wan et al. (2015; 2016; 2016; 2017) Gao et al. (2016)	Partial	No	Partial effects of turbine on platform
Geometricall y Froude scaled model	Shin (2011) Martin (2011) Koo et al. (2014) Duan (2017)	Yes	No	Turbine on platform
Redesigned blades	Martin et al. (2014) Fowler et al. (2013) Duan (2017) Du et al. (2016) Schunemann et al. (2018) Wen et al. (2020)	Yes	No	Turbine on platform
	Bayati et al. (2016; 2017; 2017)		Partially before rated wind speed	Partially both way before rated wind speed
Real-time hybrid model	Chabaud et al. (2013) Azcona et al. (2014) Sandner et al. (2015) Sauder et al. (2016) Hall et al. (2018)	Yes	No	Turbine on platform

To address the issues with the above two methods, a scaled turbine model with geometrically similar blades was used by Shin (2011), Martin (2011), Koo et al. (2014) and Duan (2017). In order to achieve the desired thrust, wind speeds were increased during the model tests (Martin, 2011; Duan, 2017). It was found that the geometrically Froude scaled model can achieve the desired thrust, but the torque was much smaller than desired. In addition, the increased wind speeds can result in extra aerodynamic loads on the tower and floating platform.

69 In order to achieve the desired thrust without increasing wind speeds considerably, low 70 Reynold's blades (or geometrically distorted blades) were adopted to replace the geometrically 71 similar blades in the model test. Martin et al. (2014) redesigned the turbine blades with the low 72 Reynolds airfoil called Drela AG04. Their results showed that the model scale thrust was close 73 to the full scale thrust, but the power and the torque coefficients in model scale were much smaller than those at full scale. Fowler et al. (2013) designed a turbine model with the low 74 75 Reynolds airfoil named Drela AG24, which was also used by Duan (2017). With the Drela series of airfoils, the chord of the redesigned blades was enlarged excessively, which resulted 76 77 in heavier blades. Low Reynolds airfoils such as NACA 4412 (Du et al., 2016), SD2030 (Schunemann et al., 2018; Wen et al., 2020) and SD7032 (Bayati et al., 2016; 2017; 2017) were 78 79 considered instead. Bayati et al. (2016; 2017; 2017) designed a model using SD7032 for DTU

10MW wind turbine (Bak et al., 2013). In their redesign, they ensured that the thrust was correctly scaled, but not the torque. Nevertheless, their test results showed that the thrust was satisfactory in all test conditions, and that the torque was also approximately modelled for wind speeds less than the rated speed. As their model was not designed to achieve torque similarity, not surprisingly, the torques corresponding to the wind speeds larger than rated speed were considerably smaller than the desired values.

86 Another approach, called Real-time hybrid model (RTHM), has also been developed. In this 87 approach, the wind turbine is replaced by a virtual subsystem and the desired aerodynamic 88 thrust is realized by a feedback controlled ducted fan or actuator (Chabaud et al., 2013; Azcona 89 et al., 2014; Sandner et al., 2015) on the top of tower. Later, Sauder et al. (2016) replaced the 90 ducted fan with a redundant cable-based actuation system to obtain the transient simulated force 91 with considering the effects of wind, wave and blade pitch control. This approach offers higher 92 bandwidth and more complex loading capabilities, however, faces a great challenge in terms of 93 complexity of equipment and controls. Then, Hall et al. (2018) used the actuation system (winch 94 and cables, pulling fore and aft on the nacelle) and combined it with a numerical wind turbine 95 tool to simulate the thrust.

In the FOWT system, the aerodynamics of turbine significantly affects the motions of 96 97 floating platforms. On the other hand, the motions of floating platforms can also affect the 98 aerodynamics (including thrust and torque) of turbines. The interaction between them plays an 99 important role in determining the dynamic behaviors of the whole FOWT systems. In the state-100 of-art studies, most approaches of model tests can only model turbine's effects (i.e., thrust) to 101 some extent. Only the approach proposed by Bayati et al. (2016; 2017; 2017) can approximately 102 model the interaction at wind speeds smaller than the rated speed. To the best knowledge of the 103 authors, no model-test method has been suggested to correctly and simultaneously model the 104 scaled thrust and torque, and so the interaction between turbine and platform based on the 105 Froude similarity law.

This paper will propose a scaled model test method, which can achieve the similarity of both thrust and torque in the whole range of wind speeds. This will establish a foundation for realizing the correct modelling of the interaction between turbine and platform during the tests of FOWT models in wind-wave basins. For the ease of discussions, a scale model of DTU 100 MW (Bak et al., 2013) wind turbine with scale factor 1:100 is designed by the present method. The proposed method is verified by the numerical and experimental results.

#### 112 2. Description of wind turbine model-test method

In the model test of FOWT systems in wind-wave basins, the interaction between turbine and floating platform mainly happens in two ways. One is the influence of turbine's aerodynamic loads (mainly the thrust) on the motion of platform, the other is the influence of platform's motion on the aerodynamic loads (mainly the thrust and torque) of turbine. To make it possible to study the full interaction by physical testing in wind-wave basins, the thrust and torque of turbine should be well scaled firstly.

113

To achieve the similarity of scaled thrust and torque simultaneously in model tests, the thrustand torque coefficients of the model and its full-scale counterpart must be the same

122 
$$C_T = \frac{T_F}{0.5\rho V_{w-F}^2 S_F} = \frac{T_M}{0.5\rho V_{w-M}^2 S_M}$$
(1)

123 
$$C_{Q} = \frac{Q_{F}}{0.5\rho V_{W-F}^{2} S_{F} R_{F}} = \frac{Q_{M}}{0.5\rho V_{W-M}^{2} S_{M} R_{M}}$$
(2)

124 in which,  $\rho$  is the density of air,  $V_w$  is the wind speed, *S* and *R* are the area and radius of rotor's 125 rotational annular plane, respectively. The subscripts *F* and *M* denote the full scale and model 126 scale, respectively.

127 The aerodynamic performance of wind turbines can be affected by many factors. These 128 include the geometrical parameters of blade, such as of the blade chord (*c*), twist angle ( $\beta$ ) and 129 relative thickness (*t/c*) distribution along the blade span. The other factors include the 130 operational parameters, such as the blade pitch angle ( $\theta$ ) and the rotational speed ( $\omega$ ).

131 Generally, a model test of FOWT system is designed based on the Froude similarity, which 132 does not satisfy the Reynolds similarity. Therefore, in the FOWT model tests, the Froude scaled 133 wind turbine may not provide correct thrust and torque to the floating platform. One approach 134 to overcome this inconsistency was proposed by Bayati et al. (2016) as indicated before. In their 135 approach, the low Reynolds blade was used which is geometrically unsimilar to the full-scale 136 blades. Main assumptions in their method were: (1) the inflow angle  $\phi$  of the model is the 137 same as that of the full-scale counterpart; (2) the blade pitch angle  $\theta$  is also the same in full 138 scale and model scale; (3) the drag force on the blade in model scale is ignored when designing 139 the model. The results by Bayati et al. (2016) showed that the model scale wind turbine with the low Reynolds blade can achieve the correctly-scaled thrust. However, it is difficult to 140 141 achieve the scaled torque in the high wind speed region larger than the rated wind speed. In this 142 study, we will propose an approach which can address this issue. Compared to the method of 143 Bayati et al. (2016), the proposed approach assumes that the twist angle in the model scale, 144 instead of pitch angle, is the same as that in the full scale. This assumption will make it easy to

- 145 determine the blade geometrical characteristics, and to obtain the lift and drag coefficients more
- accurately. In addition, this new approach will consider both of lift and drag forces on the blade,
- 147 which help achieving the correctly-scaled thrust and toque simultaneously.

#### 148 2.1 Method for designing the geometry of low Reynolds blade

As mentioned before, a low Reynolds blade for the model scale needs to be designed in FOWT model tests. The key issue of low Reynolds blade design is to determine the blade's geometrical parameters, such as chord ( $c_M$ ), twist angle ( $\beta_M$ ) and relative thickness (t/c) distribution along the blade span. As indicated before, the model twist angle is taken to be the same as the full scale. The relative thickness (t/c) distribution is determined in the same way as Bayati et al. (2016). The following discussion is presented on how to determine the chord length of the model blades.

156 In the present method, similar to Bayati et al. (2016), the inflow angle is assumed to be the 157 same in model and full scales, i.e.

158

### $\phi_M = \phi_F \tag{3}$

159 Differently from the reference (Bayati et al., 2016), we assume that the twist angle ( $\beta$ ), 160 instead of pitch angle ( $\theta$ ), is the same in the model and full scales,

161

$$\beta_M = \beta_F \tag{4}$$

162 To determine the chord length  $(c_M)$  along the blade span, we will employ a two-step approach. 163 In the first step, the preliminary chord  $(c_{M0})$  is obtained directly by matching the scaled lift force 164 along the blade span.

165 Considering the airfoil section in Fig. 2, the angle of attack  $\alpha$  is given by

166  $\alpha = \phi - (\beta + \theta) \tag{5}$ 

167 Based on the airfoil theory, the lift and drag acting on the airfoil section are

$$L = 0.5\rho c dr V_{rel}^2 \cdot Cl(\alpha)$$
(6)

169 
$$D = 0.5\rho c dr V_{rel}^2 \cdot C d(\alpha)$$
(7)

170 where,  $\rho$  is the density of air, dr is the span of the blade element,  $Cl(\alpha)$  and  $Cd(\alpha)$  are the lift 171 and drag coefficients, respectively. In the linear region (i.e. small angle of attack) of an airfoil 172 lift curve, the lift coefficient can be approximated as



175 
$$Cl(\alpha) = kl \cdot \alpha + Cl^{0} = kl \cdot (\phi - \beta - \theta) + Cl^{0}$$
(8)

176 in which, kl is the slope of airfoil lift curve in the linear region and  $Cl^0$  is the lift coefficient 177 when the angle of attack is zero.

According to the similarity theory, the relationship of lift forces between full scale ( $L_F$ ) and model scale ( $L_M$ ) is

180

$$L_F = \lambda_L^2 \lambda_V^2 \cdot L_M \tag{9}$$

181 where  $\lambda_L$  is the dimensional scale factor, equal to the ratio of full scale length to model scale 182 length;  $\lambda_V$  is the velocity scale factor, equal to the ratio of velocity of full scale to that of model 183 scale. If the Froude law is followed,  $\lambda_V = \sqrt{\lambda_L}$ . However, in practice,  $\lambda_V$  is determined, not by 184 following the Froude law, but by considering the wind-speed generation capability in laboratory 185 and by making the test cases with Reynolds number closer to that of full scale, as suggested by 186 Bayati et al. (2016).

187 Substituting Eqs. (3-8) into Eq. (9), one gets

188 
$$c_F[kl_F(\phi - \beta - \theta_F) + Cl_F^0] = \lambda_L \cdot c_{M0}[kl_M(\phi - \beta - \theta_M) + Cl_M^0]$$
(10)

189 Taking the derivative of Eq. (10) with respect to  $\beta$ , we obtain

$$c_F \cdot k l_F = \lambda_L \cdot c_{M0} \cdot k l_M \tag{11}$$

191 Thus, the preliminary model scale airfoil chord of the blade element is given by

192 
$$c_{M0} = \frac{c_F}{\lambda_L} \cdot \frac{kl_F}{kl_M}$$
(12)

which is the same as used by Bayati et al. (2016), though the twist angle is assumed to be same
for full scale and model scale in this study, while Bayati et al. (2016) assumed that the pitch
angle is the same as indicated before.

As sketched in Fig. 2, the thrust 
$$(dT)$$
 and torque  $(dQ)$  of a blade element can be estimated by

197 
$$dT = L\cos\phi + D\sin\phi \tag{13}$$

198 
$$dQ = (L\sin\phi - D\cos\phi) \cdot r \tag{14}$$

According to Eqs. (13-14), it is clear that the correctly-scaled thrust and torque can be achieved simultaneously when the airfoil lift-drag ratio and inflow angle are the same in full scale and model scale. However, the airfoil lift-drag ratio is generally different in full scale and model scale, since the low Reynolds blade is used in model tests. To make better matching of thrust and torque, a correction factor  $\gamma$  is introduced to Eq. (12) in the second step, as follows:

204 
$$c_M = \gamma \cdot c_{M0} = \gamma \cdot \frac{c_F}{\lambda_L} \cdot \frac{kl_F}{kl_M}.$$
 (15)

205 In order to evaluate  $\gamma$ , the torques in full scale and model scale are considered

$$206 \qquad dQ_F = (L_F \sin \phi - D_F \cos \phi) \cdot r_F = 0.5 \rho c_F V_{rel-F}^2 [Cl_F(\alpha_F) \sin \phi - Cd_F(\alpha_F) \cos \phi] r_F dr_F \qquad (16)$$

207 
$$dQ_{M} = (L_{M} \sin \phi - D_{M} \cos \phi) \cdot r_{M} = 0.5 \rho \gamma c_{M0} V_{rel-M}^{2} [Cl_{M}(\alpha_{M}) \sin \phi - Cd_{M}(\alpha_{M}) \cos \phi] r_{M} dr_{M}.$$
(17)

208 To fulfill the similarity in torque, they need to satisfy

$$dQ_F = \lambda_L^3 \lambda_V^2 \cdot dQ_M \tag{18}$$

210 Substituting Eqs. (16-17) into Eq. (18), the correction factor can be determined as

211 
$$\gamma = \frac{\tan \phi - 1/K_F(\alpha_F)}{\tan \phi - 1/K_M(\alpha_M)} = \frac{\tan \phi - Cd_F(\alpha_F)/Cl_F(\alpha_F)}{\tan \phi - Cd_M(\alpha_M)/Cl_M(\alpha_M)}.$$
 (19)

For estimating the value of  $\gamma$  in a simple way,  $\alpha_F$  is taken to be the attack angle at the blade tip corresponding to the rated wind speed, and  $\phi$  is taken to be the inflow angle at the tip of blade, ignoring the induced velocities, i.e.

 $\phi \approx \operatorname{arc} \operatorname{cot}(TSR_D), \qquad (20)$ 

where the  $TSR_D$  is the tip-speed ratio at design condition of the full-scale turbine.

217 According to Eqs. (10-11)

218 
$$\alpha_M = \alpha_F + \frac{C l_F^0}{k l_F} - \frac{C l_M^0}{k l_M}.$$
 (21)

219 Substituting Eqs. (20-21) into Eq. (19), one gets

220 
$$\gamma = \frac{\frac{1}{TSR_{D}} - \frac{Cd_{F}(\alpha_{F})}{Cl_{F}(\alpha_{F})}}{\frac{1}{TSR_{D}} - \frac{Cd_{M}(\alpha_{F} + Cl_{F}^{0} / kl_{F} - Cl_{M}^{0} / kl_{M})}{Cl_{M}(\alpha_{F} + Cl_{F}^{0} / kl_{F} - Cl_{M}^{0} / kl_{M})}.$$
 (22)

221 One benefit of the present method is that the re-designed low Reynolds blade does not depend

on the pitch angle which varies with wind speeds and can be applied in various model testconditions.

#### 224 2.2 Method for adjusting the pitch angle and rotational speed

In order to produce correct thrust and torque simultaneously in laboratory tests using the turbine model discussed above, the blade pitch angle and rotational speed must be properly adjusted. This section will describe the method for determining the pitch angle and rotational speed.

#### As well known, the relative axial wind speed and rotational speed can be given by

230 
$$V_{in} = (1-a)V_{w-M}$$
  $V_r = (1+b)r_M \cdot \omega_M$  (23)

where *a* and *b* are the induced factor of axial velocity and tangential velocity, respectively.

The resulting inflow velocity and inflow angle will be expressed by

233 
$$V_{rel-M} = \sqrt{\left[(1-a)^2 V_{w-M}^2 + (1+b)^2 (r_M \cdot \omega_M)^2\right]}$$
(24)

$$\phi_M = \operatorname{ar}\operatorname{cot}\frac{(1+b)r_M \cdot \omega_M}{(1-a)V_{w-M}}$$
(25)

Substituting Eqs. (24-25) in Eqs.(13-14), one gets the thrust and torque components on a bladeelement as

237 
$$dT_{M} = 0.5\rho c_{M} V_{rel-M} \cdot [(1+b)\omega_{M} r_{M} \cdot Cl_{M}(\alpha_{M}) + (1-a)V_{w-M} \cdot Cd_{M}(\alpha_{M})] \cdot dr_{M}$$
(26)

238 
$$dQ_{M} = 0.5\rho c_{M} V_{rel-M} \cdot [(1-a)V_{W-M} \cdot Cl_{M}(\alpha_{M}) - (1+b)\omega_{M} r_{M} Cd_{M}(\alpha_{M})] \cdot r_{M} \cdot dr_{M}$$
(27)

It is noted that the relationship between lift coefficient and attack angle used in Eqs. (26-27) is the real function, instead of the approximated linear relationship. And the more accurate the relationship between the attack angle and the coefficients is, the more accurate the thrust and torque will be.

In order to evaluate the induced velocity factors, the momentum theory will be applied. Fora model with *B* blades, the thrust on a blade element based on momentum equation is

245 
$$dT = \frac{1}{2}\rho S_M V_{W-M}^2 4a(1-a) = \frac{1}{2}\rho (2\pi r_M dr_M) V_{W-M}^2 4a(1-a) = B \cdot dT_M$$
(28)

246 where  $S_M = 2\pi r_M \cdot dr_M$  is the area of rotational annular plane.

247 Similarly, the torque of a blade element is

234

248 
$$dQ = (2\pi r_M dr_M \rho V_{in}) V'_r r_M = 4\pi \rho \omega_M b(1-a) V_{w-M} r_M^3 dr_M = B \cdot dQ_M$$
(29)

249 in which,  $V'_r = 2b\omega_M r_M$  is the induced tangential velocity.

250 Combining Eqs. (26-29), one obtains

251 
$$\frac{a}{1-a} = \frac{\sigma[Cl_M(\alpha_M)\cos\phi_M + Cd_M(\alpha_M)\sin\phi_M]}{4\sin^2\phi_M}$$
(30)

252 
$$\frac{b}{1+b} = \frac{\sigma[Cl_M(\alpha_M)\sin\phi_M - Cd_M(\alpha_M)\cos\phi_M]}{4\sin\phi_M\cos\phi_M}$$
(31)

253 where  $\sigma = \frac{Bc_M}{2\pi r_M}$  is the turbine solidity.

254 When the Prandtl tip loss correction factor is considered, Eqs. (30-31) can be rewritten as

255 
$$\frac{a}{1-a} = \frac{\sigma[Cl_M(\alpha_M)\cos\phi_M + Cd_M(\alpha_M)\sin\phi_M]}{4f\sin^2\phi_M}$$
(30-1)

256 
$$\frac{b}{1+b} = \frac{\sigma[Cl_M(\alpha_M)\sin\phi_M - Cd_M(\alpha_M)\cos\phi_M]}{4f\sin\phi_M\cos\phi_M}$$
(31-1)

257 in which, 
$$f = \frac{2}{\pi} ar \cos[\exp(-\frac{B}{2} \frac{R_M - r_M}{r_M \sin \phi_M})]$$

The total thrust and torque of the scaled turbine model can be calculated from Eqs. (26-27)by

260 
$$T_{M} = B \cdot \int_{0}^{R_{M}} dT_{M} = 0.5B\rho \cdot \int_{0}^{R_{M}} c_{M} V_{rel-M} \cdot [(1+b)\omega_{M} r_{M} Cl_{M}(\alpha_{M}) + (1-a)V_{w-M} Cd_{M}(\alpha_{M})]dr_{M}$$
(32)

261 
$$Q_{M} = B \cdot \int_{0}^{R_{M}} dQ_{M} = 0.5B\rho \cdot \int_{0}^{R_{M}} c_{M} V_{rel-M} \cdot [(1-a)V_{w-M} \cdot Cl_{M}(\alpha_{M}) - (1+b)\omega_{M}r_{M}Cd_{M}(\alpha_{M})] \cdot r_{M} \cdot dr_{M}$$
(33)

262 It indicates that the thrust and torque are functions of  $\theta_M$  and  $\omega_M$ .

According to Eqs. (1-2), one has

$$T_M = T_F / \lambda_L^2 \lambda_V^2, \tag{34}$$

265

264

$$Q_M = Q_F / \lambda_L^3 \lambda_V^2 \,. \tag{35}$$

Solving them together, the pitch angle  $\theta_M$  and rotational speed  $\omega_M$  can be determined for each wind speed. In the equations,  $T_F$  and  $Q_F$  are assumed to be given. If that would not be the case, one can estimate them using numerical tools.

The model test method for achieving similarity of both thrust and torque of model- and fullscale wind turbines is summarized in Fig. 3. It will be validated with numerical and experimental results in the following section.





#### Fig. 3. Wind turbine model-test method

The main features of the present method and its difference from that presented by Bayati et 274 al. (2016) are summarized in Table 2. There is one point to be further discussed here, which is 275 about the rotational speed. In the present method, the rotational speed is determined using Eqs. 276 (34-35) to ensure the similarity of both thrust and torque while it was determined by  $\omega_F \lambda_V / \lambda_L$ 277 in Bayati et al. (2016). The Froude scaling of rotational speed should be  $\omega_F / \sqrt{\lambda_L}$ . The results 278 below will show that the rotational speed from Eqs. (34-35) is in the range of 279  $\omega_F / \sqrt{\lambda_L} < \omega_M < \omega_F \lambda_V / \lambda_L$  at least for the cases considered. The choice of  $\omega_F \lambda_V / \lambda_L$  by 280 Bayati et al. (2016) ensures that the Tip Speed Ratio (TSR) is the same for both model and full 281 282 scales, which is not the case in the present method. Certainly, a mismatching in the TSR together with other parameters such as the Froude number for air, blade geometry and so on 283 284 will lead to inconsistence of the detailed aerodynamics between model- and full-scale turbines. Nevertheless, the present method realizes the similarity of both thrust and torque, which are 285 286 important for studying the interaction between turbine and floating platform by physical tests 287 in wind-wave basin.

- 288 Table 2
- 289 Comparison of main features of different scaling methods

Main Parameters of scaled model	Method in Bayati et al. (2016)	Present method
Thrust	correctly scaled	correctly scaled
Torque	incorrectly scaled at high speed	correctly scaled
Tip Speed Ratio	correctly scaled	incorrectly scaled
Twist angle	incorrectly scaled	correctly scaled
Pitch angle	correctly scaled	incorrectly scaled
Rotational speed	incorrectly scaled	incorrectly scaled
Blade section	incorrectly scaled	incorrectly scaled

Turbine diameter	correctly scaled	correctly scaled

292

#### 3. Numerical validation of wind turbine model-test method

The present method is firstly validated by using the commercial CFD software Star-CCM+ 293 294 (Siemens, 2017). For ease of discussions, a scaled model of DTU 10MW (Bak et al., 2013) 295 wind turbine with scale factor 1:100 is re-designed by the present method. Then, the thrust and 296 torque with different velocity scale factors are calculated. In preparation for the model test of 297 FOWT system in wind-wave basins, the wind speeds are selected by considering the wind-298 speed generation capability in wave basins.

299



301

#### 3.1 Design of model scale wind turbine 302

303 In this study, the Selig-Donovan low Reynolds airfoil SD7032 following Bayati et al. (2016) is employed and the corresponding Reynolds number is about  $12.5 \times 10^4$ . Comparisons of the 304 305 airfoil lift and drag coefficients between model- and full-scale airfoils are shown in Fig. 4. It 306 can be seen that the airfoil in full scale and model scale blades have obviously different stall angles, which are about 15° and 11°. In standard working condition, most length of full scale 307 308 blades works far from stall and the attack angle is about in the range of [-2, 9]. For this range, 309 the lift coefficients of airfoil in full scale and model scale blades are approximately linear, and 310 the slope of the former is larger than the latter. According to Eq. (12), the chord of preliminary 311 model scale airfoil chord of the blade element will be larger than the geometric scaled one. 312 According to Eq. (8), the lift curves can be approximated as a linear relationship to angle of

313 attack

314 
$$Cl_F(\alpha_F) = kl_F \cdot \alpha_F + Cl_F^0 = 0.1201 \cdot \alpha_F + 0.34233$$
 (36)

315 
$$Cl_{M}(\alpha_{M}) = kl_{M} \cdot \alpha_{M} + Cl_{M}^{0} = 0.08766 \cdot \alpha_{M} + 0.42751$$
 (37)

316 According to Eq. (12), the preliminary model scale airfoil chord is

317

$$c_{M0} = \frac{c_F}{\lambda_L} \cdot \frac{kl_F}{kl_M} = 1.37 \cdot c_F / \lambda_L \tag{38}$$

With the enlarged factor (Eq.22) at the rated wind speed equal to 1.0836, one obtains

$$c_M = \gamma \cdot c_{M0} = 1.4845 \cdot c_F / \lambda_L \tag{39}$$

The relative thickness distribution is determined in the same way as Bayati et al. (2016). With using the same twist angle as that of full scale blades, the model scale blade can be obtained, as shown in Fig. 5.







**Fig. 7.** Pitch angle and rotational speed for different wind speeds ( $\lambda_V = 3$ )

The pitch angle and rational speed can be determined from Eqs. (34-35) as mentioned before. Fig.6 and Fig.7 show the results for  $\lambda_v = 2$  and  $\lambda_v = 3$ , respectively. In the figures, the green solid lines present the results obtained with the method proposed by Bayati et al. (2016), the black solid line is obtained with direct Froude scaling, while the red dash-dotted lines and blue dashed lines are obtained with the present method. This clearly shows the rotational speeds of the present method is in between that used by the Bayati et al. (2016) and that based on the full Froude scaling.

#### 339 *3.2 Computational set-up and validation*

In order to ensure that the results of Star-CCM+ are credible, the aerodynamic forces of full scale DTU 10MW wind turbine are calculated and compared with the reference results firstly. The computational domain setup is shown in Fig. 8. The maximum mesh size is about 0.23*R*, and the mesh is encrypted to a maximum size of 0.057R near the wind turbine. The maximum and minimum size on the surface of blade is about 0.0018R and 0.00022R respectively, 2 layers of boundary layer mesh are generated with a total layer thickness of 0.0007R and a progression factor of 1.2. The total number of cells in the computational domain is about  $1.73 \times 10^{6}$ .





Fig. 9. Thrust and torque of full scale DTU-10MW wind turbine for different wind speeds

354 The thrust and torque of full scale DTU-10MW wind turbine against wind speed are shown 355 in Fig. 9. In this figure, the green solid lines present the results obtained by HAWCStab2 (Bak 356 et al., 2013), while the red points are the results obtained by Star-CCM+ in this study. It should 357 be noted that there are two points when the wind speed is 11.4m/s, which is the rated wind speed of full scale DTU-10MW wind turbine. One point is the results with the rotational speed 358 359 of 9.16rpm, and the other is the results with the rotational speed of 9.6rpm. In fact, the rated wind speed is the critical state of full scale DTU-10MW wind turbine. As known, the TSR (7.5) 360 361 is kept when the wind speed is less than the rated one, and the maximum power coefficient is used to capture the wind energy. By this way, the rotational speed at the rated wind speed should 362 be 9.16rpm, the thrust and torque will be the smaller and larger values respectively. However, 363 364 the maximum rotational speed (9.6rpm) is kept when the wind speed is larger than the rated one, and the rated power is used to capture the wind energy. By this way, the rotational speed 365

at the rated wind speed should be 9.6rpm, and the thrust and torque will be the larger and smaller 366 values respectively. From the figure, it can be seen that the results of Star-CCM+ agree well 367 368 with the results of HAWCStab2 which is based on BEM, the errors for all considered wind 369 speeds do not exceed 6.5%. Fig. 9 indicates that the results of Star-CCM+ are credible.







377 Fig. 10 presents the thrust and torque of the redesigned model scale wind turbine with velocity scale factor 2. The green solid lines are the directly scaled results obtained by Eqs. (34-378 379 35), the blue crosses are obtained by the present design method with the airfoil chord  $c_{M0}$ defined in Eq.(38), and the red points are obtained by the present design method with the airfoil 380

chord  $c_M$  defined in Eq.(39). With the scaled pitch angle and rotational speed, shown in Fig. 6, 381 382 the scaled thrust and torque can be achieved simultaneously in the whole range of wind speeds. 383 It should be noted that, in low wind speed region (i.e. lower than rated wind speed), the result 384 with  $c_M$  is in better agreement with the directly scaled result (i.e. the green solid line) than that 385 with  $c_{M0}$ . As we know, in the low wind speed region, the angle of attack is commonly close to 386 the one with the maximum lift coefficient. To achieve enough thrust and torque, the airfoil chord should be slightly larger since the lift-drag ratio of the redesigned low Reynolds blade is 387 relatively low. This is the possible reason why the result with  $c_M$  shows to be better than that 388 389 with *c*<sub>M0</sub>.

A scaled wind turbine with velocity scale 3 is also carried out to check whether the redesigned blade obtained by the present method is suitable for cases with different velocity scale factor. It is clear that the scaled thrust and torque can also be achieved simultaneously in the case with velocity scale 3, as shown in Fig. 11. It indicates that the redesigned blades obtained by the present method can be used for cases with different velocity scale factors. The only thing we should do is to adjust the pitch angle and rotational speed for different cases, as shown in Fig. 7.



Fig. 12. Physical model of redesigned blade



Fig. 13. Experimental setup

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#### 4. Experimental validation of wind turbine model-test method

For the model test of wind turbines in wind tunnel, the increased wind and rotational speeds are generally used to maintain the Reynolds scaling. In this paper, the wind tunnel test is to further validate the present method and preparing for the model test of FOWT system in wind-wave basins. Therefore, the wind speeds which can be realizable by the wind generating system in wave basin are employed. And the velocity scale factor in wind tunnel test is 2.

The model test is carried out in a boundary layer wind tunnel, which has the dimensions of  $18m \times 4m \times 3m$ . The physical model of redesigned low Reynolds blade is shown in Fig. 12, and the experimental setup is shown in Fig. 13. In experiment, the thrust is measured by a 6-component balance under the bottom of tower and the torque is measured by a contactless torque sensor on rotation axis.







Wind speed (m/s)

Fig. 15. Experimental-numerical errors of thrust and torque relative to scaled results

Wind speed (m/s)

The experimental results when the velocity scale factor equals to 2 are shown in Fig. 14. In the figure, the green solid lines are the directly scaled results obtained by Eqs. (34-35). They

422 are also the desired scale thrust and torque of model scale wind turbine for different wind speeds. 423 The operating conditions in experiment are also solved based on the green solid lines. The red 424 points are the results of experiment and the black dashed lines are obtained with Star-CCM+ 425 by using the operating condition in experiment. It can be seen that the results of Star-CCM+ 426 agree well with the experimental results for most of the wind speeds, which further indicates 427 that the results of Star-CCM+ are credible.

Fig. 15 shows the error relative to the desired scale thrust and torque (green lines in Fig. 14).
It can be seen that the results of Star-CCM+ agree very well with the desired scale results, and
most of the errors are less than 4%. The numerical results indicate that the operating conditions
in experiment are credible and they can make the model scale wind turbine achieve the desired
scale thrust and torque simultaneously.

433 When the wind speed is less than the rated wind speed (5.7 m/s), the experimental results 434 have the same trends as the desired scale results, obtained with Eq. (34-35). Although the 435 experimental thrust force is smaller than the scaled value, most of the errors are less than 10% 436 which are acceptable. In addition, the scaled torque is also achieved simultaneously. When the 437 wind speed is larger than the rated wind speed, the experimental thrust force agrees well with 438 the scaled value. And the experimental torque is smaller than the scaled value, however, most 439 of the errors are less than 5%. In summary, although the thrust and torque do not agree 440 excellently with the desired scale results, the corresponding errors can be considered as 441 acceptable. To the best of our knowledge, no other model test methods can achieve similarity 442 of both thrust and torque at high wind speed, greater than the rated wind speeds.

443 From Fig. 14, the only major discrepancies in between the results of experiment and Star-CCM+ correspond to the wind speed conditions of 9m/s~11m/s. This may be caused by the 444 445 variation of relative inflow velocity. In Star-CCM+, the wind speed is uniform and the wind 446 turbine is rigid. While in physical experiment, the wind speed is not necessarily uniform and 447 completely stable. Besides, the 6-component balance under the tower is a flexible structure, 448 which makes the tower vibrate obviously with the increasing of wind speed. As a result, the 449 attack angles of blade elements will oscillate around its equilibrium position. Under the wind 450 speed conditions of 9m/s~11m/s, the attack angles of more than 50% of blade span are negative. 451 According to the lift and drag curves of airfoil used in model scale blades (Fig. 4), the vibration 452 of attack angle will lead to the decreasing and increasing of lift and drag component 453 respectively. According to Eq. (13) and Eq. (14), the torque of wind turbine will be less 454 obviously than the results of Star-CCM+.

#### 5. Conclusions and discussions

457 In this study, a wind turbine model-test method is proposed for achieving similarity of both model- and full-scale thrust and torque. This is achieved by redesigning the model blades and 458 459 by adjusting the pitch angle of the blades and rotational speed of the turbine. In the present method, the model blade is redesigned by keeping its twist angle, instead of pitch angle, same 460 461 as that of the full-scale blades, which makes it possible for the model to suite a wide range of 462 wind speeds. In addition, this new approach considers both lift and drag forces when 463 determining the chord length of the airfoil of the model blade, which helps achieving the 464 correctly-scaled thrust and toque. During model tests, the pitch angle of the blades and the 465 rotational speeds of turbine are adjusted by ensuring that overall torque and thrust are similar 466 for both model- and full-scale turbines in the full range of wind speeds up to the cut-off speed. 467 Numerical simulations and wind tunnel model tests are carried out to validate the new 468 method. In the numerical validation, two scaled wind turbines with different velocity scale 469 factors, i.e.  $\lambda_v = 2$  and  $\lambda_v = 3$ , are simulated by the CFD software Star-CCM+. Numerical 470 results show that the correctly scaled thrust and torque can be achieved simultaneously for all 471 the wind speeds under which the turbine would be operated for both the velocity scale factors. 472 In the experimental validation, a scaled wind turbine with velocity scale factor  $\lambda_v = 2$  is tested 473 in a boundary layer wind tunnel. The results demonstrate that the proposed method can realize 474 the similarity of thrust and torque simultaneously for all the wind range of speeds.

To the best knowledge of the authors, this is the first time to report such model test method which can achieve similarity of both model- and full-scale thrust and torque for a wide range of wind speeds. This lays a sound foundation for achieving correct interaction during the model tests of floating offshore wind turbines in wind-wave basins. The method will be applied to such tests in our future work.

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