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Citation: Jebelli, M., Shariloo, K. & Masdari, M. (2024). Vortex-plate interaction in wake-induced vibration: A study of single and parallel flat plates behind a cylinder. *Ocean Engineering*, 300, 117340. doi: 10.1016/j.oceaneng.2024.117340

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Vortex-Plate Interaction in Wake-Induced Vibration: A Study of Single and Parallel Flat Plates Behind a Cylinder

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1 Abstract

2 This computational research investigates the oscillations caused by the wake of a circular
3 cylinder on either a single or dual flat plates in a laminar flow scenario. The findings emphasize
4 the crucial role of the flat plate placement, which can significantly impact their ability to serve
5 as control devices or experience noticeable oscillations. When a plate is positioned beyond
6 the crucial juncture, it causes significant amplitude oscillations. These oscillations decrease
7 when the plate is moved further away from the horizontal axis. Furthermore, at specific short
8 horizontal distances, there are noticeable oscillations with a moderate amplitude. These
9 oscillations are caused by changes in the immediate wake structure. Oscillations in dual
10 plates, which can oscillate independently, occur only when there is sufficient room for shear-
11 layer contact. By adjusting the vertical distance, we may see different effects on vortex
12 damping and plate oscillation. Increasing the gap size reduces the area where there is no
13 oscillation. This experiment also encompasses a range of lower velocities, demonstrating that
14 although the plate(s) reliably replicate the shedding frequency of the upstream cylinder in all
15 configurations, this frequency can be affected and altered in tiny gaps. An analysis of single
16 and dual plates reveals their effectiveness in designing energy harvesting devices, with two
17 plates offering a slight advantage by utilizing separate components within a very similar area.

18 Keywords

19 Cylinder; Flat plate; Vortex shedding; Wake-Induced Vibration; Laminar flow

20 1 Introduction

21 Cylinders are a fundamental topic in fluid mechanics engineering for studying the flow
22 dynamics around bluff bodies. Vortex formation in the wake zone is a prominent feature of
23 the flow around a circular cylinder. This phenomenon has been intensively studied due to its
24 relevance to various applications. These encompass the dynamic interactions between fluids
25 and structures, specifically in the context of buildings like bridges, skyscrapers, and wind
26 turbine blades (Jain et al., 1996; Yang et al., 2004; Arrigan et al., 2011). Additionally, research
27 has been conducted on the phenomena of noise and thermal transmission (Zhang and Øiseth,
28 2021; Xu et al., 2019).

29 The Reynolds number, which is determined by the diameter of a cylinder and the velocity of
30 the undisturbed flow, ($Re = \rho UD / \mu$), has a significant effect on the wake flow. It affects
31 different flow characteristics, such as the shear layers and the pattern of vortex shedding.
32 This effect becomes noticeable when the Reynolds number exceeds 47. Vortex-Induced
33 Vibration (VIV) is a self-excited oscillation that can occur when a structure interacts with
34 shedding vortices (Blevins, 1977). Wake-induced vibration (WIV) occurs when vortices are
35 discharged from a bluff body located upstream (Zdravkovich, 1988; Bokaian and Geoola,
36 1984). A multitude of reviews have been carried out on these phenomena, as well as on Flow-
37 induced motions, in the past few decades (Bearman, 1984; Griffin and Ramberg, 1982;
38 Parkinson, 1989; Sarpkaya, 1979; Williamson and Govardhan, 2004).

39 The primary aim of early research on Flow-Induced Motions (FIMs) was to get a better
40 understanding of these phenomena and develop effective strategies to mitigate their
41 negative effects. Nevertheless, over the past ten years, there has been a growing interest in
42 FIMs as a practical method for producing sustainable energy. The combination of a cylindrical

43 body and a connected or separate flat plate, known as a simple and effective passive method,
44 has been extensively researched mostly for controlling and reducing FIMs, and sometimes, to
45 enhance them for the purpose of energy harvesting.

46 Studies on the effects of a long splitter plate attached to the back of a cylinder have shown a
47 notable decrease in vortex shedding and a large reduction in drag (Roshko, 1955; Apelt et al.,
48 1973; Apelt and West, 1975). Using a separate plate has been found to be advantageous in
49 controlling vortex shedding (Unal and Rockwell, 1988). The essential thresholds for
50 influencing the flow surrounding a circular cylinder in laminar conditions have been
51 determined at horizontal separations of 2.5 and 2.6D (Lin and Wu, 1994; Hwang et al., 2003).
52 To achieve the best possible suppression, Dehkordi and Jafari discovered that the necessary
53 horizontal distance and the vertical location of the plate are key factors that interact with
54 each other. Their research revealed that parallel splitters can reduce drag force by up to 20%
55 (Dehkordi and Jafari, 2010).

56 The use of a splitter plate is anticipated to deliberately delay the shear layer interaction,
57 leading to a subsequent reduction in hydrodynamic forces and, consequently, alleviating the
58 impact of flow-induced vibrations (FIV) on a circular cylinder (Tang et al., 2023).

59 Studies have been conducted to examine the effects of incorporating a splitter plate on
60 Vortex-Induced Vibrations (VIV). These investigations have been done using experimental
61 methods (Assi and Bearman, 2015; Stappenbelt, 2010; Assi et al., 2009) as well as numerical
62 analyses (Sahu et al., 2019; Sun et al., 2020; Zhang et al., 2021). The size of the attached plate
63 significantly influences the FIV behavior of the cylinder, causing self-restrained oscillations for
64 shorter plates and uncontrolled galloping for longer plates (Wu et al., 2023). A plate that is
65 sufficiently extensive will effectively suppress oscillations (Zeng et al., 2023). In contrast, a

66 pliable plate with a length in the direction of the flow that is greater than the diameter of the
67 cylinder significantly reduces vibrations (Cui et al., 2022). Moreover, a study conducted on a
68 cylinder combined with a separate flexible plate identified two clearly defined phases of
69 reaction, distinguished by the horizontal distance between the objects. These phases were
70 classified as Vortex-Induced Vibration (VIV) for shorter gaps and galloping for longer gaps
71 (Mittal and Sharma, 2022).

72 Recently, there has been a trend in study towards developing methods for energy harvesting
73 from different setups comprising a cylinder and a plate. An area of focus is the
74 implementation of a piezoelectric plate on the back side of the cylinder, which has been
75 thoroughly investigated. An et al. introduced a novel energy conversion device called the
76 vortex-induced piezoelectric energy converter (VIPEC). This device utilizes the pressure
77 difference generated by vortices to activate a plate, compress piezoelectric patches, and
78 transform fluid dynamic energy into electrical energy. According to their research, changing
79 the length of the plate affects how it interacts with shear layers and delays the shedding of
80 vortices (An et al., 2018). Su and Lin developed a vortex-induced bi-directional piezoelectric
81 energy harvester that can efficiently capture wind energy from two perpendicular directions.
82 The device is designed in a U-shape and has the ability to modify the aspect ratio, which in
83 turn allows for the control of wind lock-in speed (Su and Lin, 2020). Wang et al. proposed
84 attaching a rectangular stationary plate behind a previous circular cylinder to enhance the
85 energy extraction from its fluid-induced vibrations (FIV) (Wang et al., 2021). In 2016,
86 Armandei and Fernandes conducted a study on the extraction of marine current energy using
87 a device called WIV. They used a modified cylinder with a flat plate that has the ability to
88 rotate. According to their research, the buffeting response was found to be influenced by the
89 position of the elastic axis and the spring rate, rather than the gap between the plate and

90 cylinder (Armandei and Fernandes, 2016). Derakhshandeh conducted an additional study
91 where they developed a self-excitation thin piezoelectric actuator that was affected by the
92 vortex shedding occurring behind a stationary cylinder. Derakhshandeh (2022) found a strong
93 correlation between the frequency of the vortex and the displacement of the tip, indicating a
94 close relationship with the maximum power generated.

95 After examining the mentioned research, it is evident that flat plates have a broad range of
96 uses that go beyond their traditional role in controlling and reducing Flow-Induced Motions
97 (FIMs). They can potentially enhance the movement of cylinders and can also be used for
98 directly harvesting energy. However, there is a lack of research on detachable plates, which
99 are particularly sensitive to geometric and flow dynamics. Most existing studies focus on
100 analyzing the flow field and regulating the Fluid-Induced Motion (FIM).

101 This study aims to elucidate two aspects by investigating the interactions between vortices
102 and plates. The plates, either single or dual, are positioned in the flow behind a cylinder
103 located upstream. Initially, the study aims to investigate the interaction between the flow
104 created by the wake and the oscillating flat plates, as well as the resulting effects on the
105 cylinder located upstream. This examination will enhance our understanding of fluid
106 mechanics and the interaction between flow configurations and structural behaviors. The
107 study primarily aims to determine certain conditions that are favorable for extracting energy
108 from the passage of air behind an object. This particular segment, which is the main focus of
109 the inquiry, has the potential to establish the foundation for the development of new energy-
110 harvesting technologies that utilize fluid dynamic forces.

111 The article commences with an introduction to the issue at hand, succeeded by an exposition
112 of the governing equations, numerical methodology, and method validation in Section 2.

113 Section 3 is bifurcated into two segments, focusing on the exposition of results and
114 discussions pertinent to the WIV of the plates. The paper culminates with a synthesis of the
115 principal findings and conclusions in Section 4.

Nomenclature

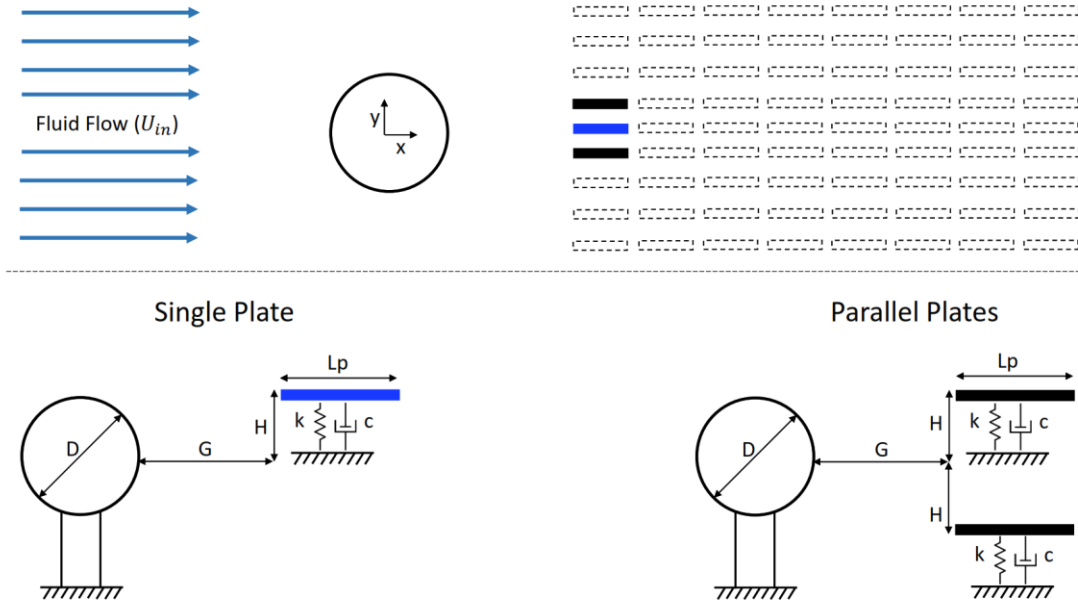
A^*	Plunging Amplitude (A/D)
D	Circular Cylinder Diameter
L_p	Flat Plate Length
G	Non-Dimensional Horizontal spacing
m^*	Mass Ratio
ζ	Damping Ratio
CL	Lift Coefficient
CL_{rms}	Root Mean Square of Lift Coefficient
CL_{mean}	Mean lift coefficient
CD_{mean}	Mean Drag Coefficient
C_p	Pressure Coefficient
C_{p_b}	Base Pressure Coefficient
ρ	Fluid Density
m	Body Mass
m_A	Added Mass
K	Transverse Stiffness Factor
U_{in}	Free Stream Velocity
U_r	Reduced Velocity
t	Physical Time
T	Non-Dimensional Time
St	Strouhal Number
VSA	Vibration per swept area

117 2 Problem description, governing equations, numerical methodology 118 and validation

119 This section begins by presenting the problem statement, outlining the model configuration,
120 and explaining the governing equations relevant to the fluid-solid interaction problem. The
121 text explains the numerical method used to solve the equations and concludes with a
122 validation of the method.

123 2.1 Problem description

124 This study examines the transverse free oscillation of a single or parallel flat plates mounted
125 in the wake of an upstream fixed circular cylinder at a Reynolds number of 100, as shown in
126 Figure 1. The study of the single plate involves examining different horizontal and vertical
127 positions, including the centerline of the wake and four other vertical positions ($H=0, 0.25,$
128 $0.5, 0.75, 1$). This analysis covers a total of 40 cases to investigate the asymmetrical
129 configurations. When considering the parallel plates, similar configurations are analyzed with
130 a minimum vertical gap of $H=0.25$, resulting in 32 variations (eight horizontal positions are
131 considered for each scenario).



132

FIG. 1) Model configuration for single and parallel flat plates

133

The plate's length (L_p) is set equal to the cylinder's diameter (D), and the incoming flow is

134

defined by a constant velocity represented as U_{in} . The device holding the plate(s) includes a

135

spring with stiffness k and a mechanical damping coefficient c , which can be set to zero to

136

allow for higher vibration amplitudes.

137

2.2 Governing Equations

138

The flow field around objects are characterised through the unsteady Navier-Stokes

139

equations, predicated on the assumption that the flow is Newtonian and incompressible. The

140

equations are rendered in a non-dimensional format, employing an Arbitrary Lagrangian-

141

Eulerian (ALE) frame of reference, as articulated in equations (1) and (2).

$$\frac{\partial u_i}{\partial x_i} = 0 \quad (1)$$

142

$$\frac{\partial u_i}{\partial t} + (u_j - \hat{u}_j) \frac{\partial u_i}{\partial x_j} = -\frac{\partial p}{\partial x_i} + \frac{1}{Re} \frac{\partial^2 u_i}{\partial x_j^2} \quad (2)$$

143 The variables x , t , and p are non-dimensionalized with respect to the cylinder's diameter and
144 the reference time ($t_{ref} = c/U_\infty$), and $\rho_\infty U_\infty^2$, respectively.

145 The transverse oscillation of the plate(s) is described as a system consisting of mass, spring,
146 and damper, as shown in equation (3):

$$m\ddot{y} + c\dot{y} + ky = F_y(t) \quad (3)$$

147 Where y , \dot{y} and \ddot{y} represent the displacement, velocity, and acceleration of the plate(s) in the
148 transverse direction, respectively. The parameter m represents the total mass that oscillates
149 per unit length, whereas c and k represent the structural damping coefficient and the
150 stiffness of the spring, respectively. $F_y(t)$ represents fluctuating transverse force per unit
151 length acting on the oscillating system.

152 The flow field is determined at each time step by solving the flow governing equations, and
153 the hydrodynamic forces on the plate are calculated accordingly. The computed forces are
154 then passed to a user-defined function, and the equation of motion is solved to determine a
155 new position for the plate. The new plate characteristics are passed back to the solver, the
156 current location is adjusted, and the grids in the dynamic mesh zone are regenerated to
157 prepare for calculating the flow field in the next step.

158 Flow-induced forces on objects are calculated by integrating the normal pressure and shear
159 stress over their surfaces. The instantaneous lift and drag coefficients are defined as $C_L(t) =$
160 $2F_y(t)/(\rho U^2 D)$ and $C_D(t) = 2F_x(t)/(\rho U^2 D)$, respectively. Here F_y and F_x denote the force
161 components in the cross-flow and inline directions, respectively. The plates' dynamic behavior
162 is defined by two parameters: vibration amplitude and frequency. The vibration amplitude is
163 made dimensionless by dividing it by the cylinder diameter, which is identical to the length of

164 the plate, and is defined as $A^* = A/D$. The vibration frequency variation is represented by
165 the frequency ratio, which is normalized by the object's natural frequency (F_{os}/F_n).

166 The main goal of the current study is to finding configurations that have a higher capacity for
167 energy harvesting. The extracted energy is directly linked to the vibration amplitude of the
168 objects. A smaller spatial footprint of the system is more advantageous for designing an
169 energy harvesting mechanism, which enhances the system's overall effectiveness. The
170 Vibration per Swept Area (VSA) statistic is introduced in (4) to enhance the comparison of
171 different configurations. This parameter offers understanding of the system's dimensions and
172 its correlation with vibration intensity.

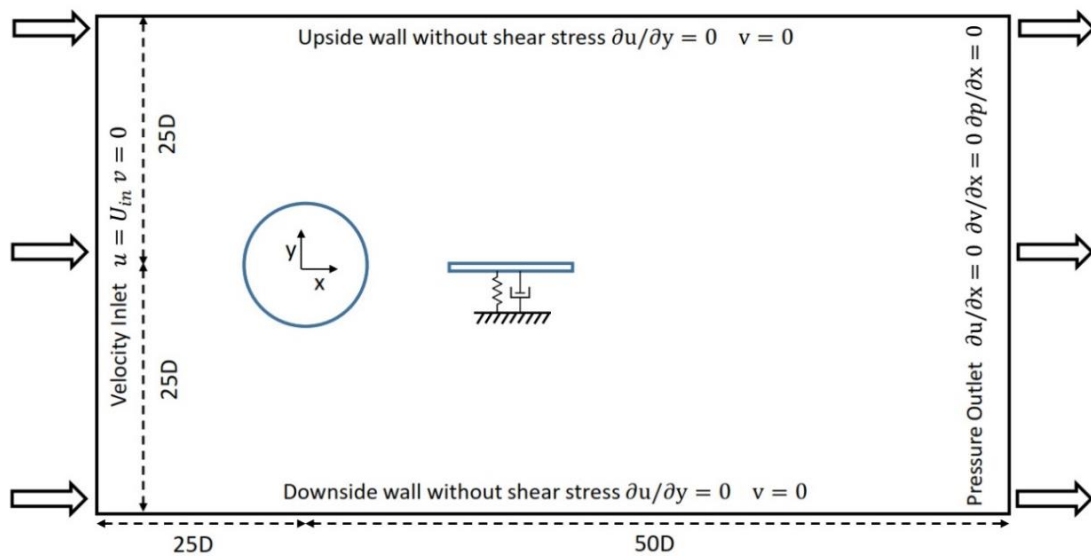
$$VSA = \frac{\sum_{i=1}^n A_i^*}{L W} \quad (4)$$

173 Where A_i^* represents the vibration amplitude of the i^{th} object, L represents the total
174 horizontal extent of the system, and W indicates the maximal vertical distance traversed by
175 the objects during a single vibration cycle.

176 2.3 Numerical Method

177 This study utilizes the finite volume method to quantify the governing equations. Time aspects
178 are divided into discrete intervals using a first-order implicit time integration method.
179 Convection terms are handled using a second-order upwind technique, while diffusion terms
180 are dealt with using a cell-based least square scheme for spatial discretization. The pressure-
181 velocity coupled equations are solved using the SIMPLE technique. In addition, a diffusion-
182 based smoothing method is used in this study coupled with a User-Defined Function created
183 in the C Programming Language to analyze the object's response to Wake-Induced Vibration
184 (WIV). The size of the computational area, especially the blockage effect, is crucial in
185 numerical simulations. Choosing a region that is too small can affect the results. Transverse

186 vibration in Flow-Induced Vibration (FIV) models might worsen inaccuracies. Studies have
 187 shown that a computational area with a blockage of more than 2.5% around a circular cylinder
 188 could cause hysteresis in the vibration response when entering the lock-in regime. Therefore,
 189 for this research, setting lateral limits with a blockage of 2% is considered appropriate.
 190 (Prasanth et al., 2006), (Prasanth and Mittal, 2008), and (Prasanth and Mittal, 2009).
 191 The boundary conditions are specified in Figure 2. The center of the cylinder is set as the origin
 192 of the coordinate system, and the rectangular computing domain spans $75D \times 50D$. The
 193 upstream and lateral limits are located at a distance of $25D$ from the center of the cylinder,
 194 while the downstream boundary is set at $50D$ from the origin to prevent computational
 195 errors.

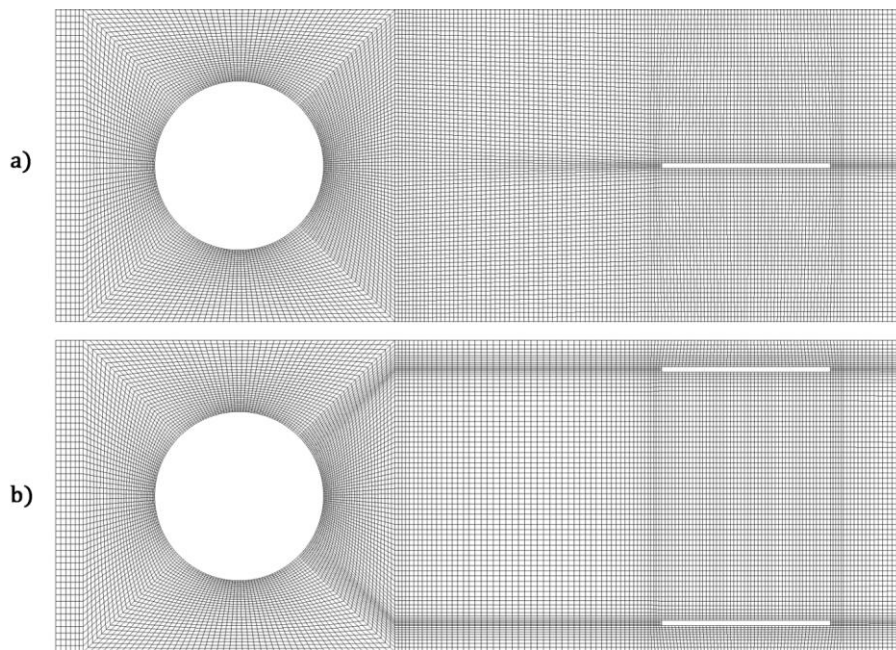


196 **FIG. 2) Computational domain and boundary conditions**

197 The inlet boundary condition is under a uniform velocity condition ($u = U_{in}, v = 0$). The
 198 velocity at the outflow has a predetermined zero normal gradient, and the pressure is set as
 199 the reference value of zero. The boundaries are defined by the normal velocity component,
 200 where the stress vector component parallel to the boundary is zero ($\partial u / \partial y = 0, v = 0$). Both
 201 the cylinder and the plate(s) have a no-slip condition applied to their surfaces.

202 **2.4 Method validation**

203 Prior to thoroughly examining the main cases, it is crucial to perform a grid independence
204 study and validate the numerical methods. Two specific situations are selected to verify the
205 accuracy of the procedure. The first scenario deals with the natural oscillation of a single flat
206 plate positioned behind a circular cylinder, whereas the second scenario features two parallel
207 flat plates mounted in the wake and immobile. Therefore, two distinct grids are built, as seen
208 in Figure 3.



209 **FIG. 3) Grid distribution. a) Single Plate b) Parallel Plates.**

210 The cylinder's surface in both grids is divided into four distinct sections: front, top, bottom,
211 and rear sides. The rear side is especially significant because of the increased flow gradients
212 in the nearby wake zone. The use of flat plates in this location increases the complexity of the
213 flow dynamics in that specific region. As a result, the grids are made more detailed in specific
214 zones and surrounding the entities, then gradually made less detailed in farther places to
215 decrease processing requirements. Three grids with different mesh densities are created for
216 each scenario to ensure grid-independent results. In both cases, the cylinder upstream does

217 not move, while the plate(s) downstream are allowed to vibrate freely in the direction of the
 218 crossflow. Tables I and II provide the root mean square lift coefficient, CL_{rms} , and mean drag
 219 coefficient, CD_{mean} , for the cylinder with maximum vibration amplitude, as well as the CL_{rms}
 220 for the plate(s).

221 **TABLE I) Grid dependency study (G=2-H=0, Single plate)**

Grid	Cylinder Nodes	Cylinder		Single Flat Plate	
		CL_{rms}	CD_{mean}	A^*_{max}	CL_{rms}
SP-1	180+70	0.069	1.18	0.28	0.267
SP-2	180+80	0.073 (5.8%)	1.194 (1.2%)	0.282 (0.7%)	0.273 (2.2%)
SP-3	180+90	0.075 (2.7%)	1.199 (0.4%)	0.281 (0.35%)	0.274 (0.3%)

222 Increasing the number of cells on the back side of the cylinder improves the accuracy of
 223 predicting hydrodynamic parameters for both the plate and the cylinder, particularly the CL_{rms}
 224 of the cylinder. The results show that the increased mesh density of grid SP-2, with 33% more
 225 nodes on the back of the cylinder along the wake, confirms its effectiveness in predicting
 226 accurately.

227 **TABLE II) Grid dependency study (G=3-H=1, parallel plates)**

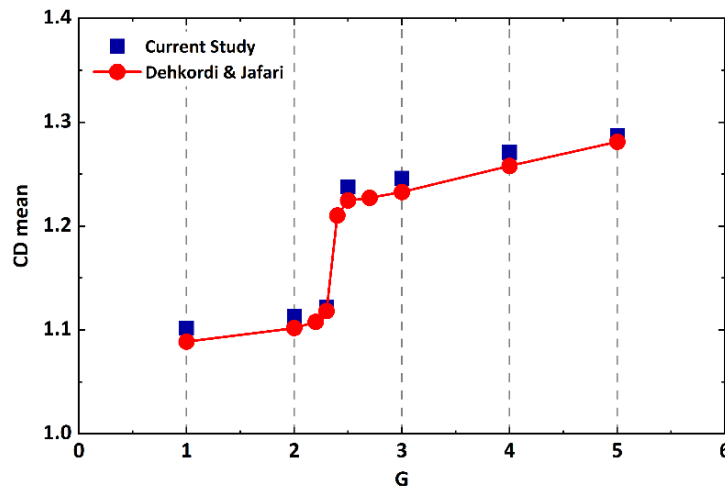
Grid	Cylinder Nodes	Cylinder		Parallel Flat Plate(s)	
		CL_{rms}	CD_{mean}	A^*_{max}	CL_{rms}
PP-1	180+80	0.24	1.297	0.356	0.38
PP-2	180+90	0.249 (3.75%)	1.30 (0.3%)	0.36	0.386
PP-3	180+100	0.253 (1.6%)	1.30 (0%)	0.36	0.389

228 When using parallel plates, results from grid PP-2 show that adding 50% extra cells to the rear
 229 side of the grid ensures that the results are not affected by the grid resolution. Both grids
 230 were analyzed at different time intervals for temporal resolution, with a non-dimensional
 231 time step of $\Delta t = 0.002$ ($t_{Non-D} = tU/D$ found to be sufficiently small to guarantee
 232 independence, as demonstrated in Table III.

Table III) Time step dependency study for single and parallel plates

Grid	Time Step	Flat Plate Amplitude	
		G=3-H=0	G=3-H=1
Single Plate	0.004	0.54	0.351
	0.002	0.515	0.32
	0.001	0.51	0.314
Parallel Plates		G=3-H=1	G=4-H=1
	0.004	0.377	0.389
	0.002	0.36	0.384
	0.001	0.35	0.382

234 As stated earlier, to establish a validated computational method requires investigating two
 235 distinct cases. The first case aims to confirm the effectiveness of the method in simulating the
 236 flow dynamics around a cylinder and a single flat plate, while the second case concentrates
 237 on parallel plates. Jebelli and Masdari (2022) examined the Flow-Induced Vibration (FIV) of a
 238 single flat plate positioned along the central axis of the wake in their previous study.
 239 Therefore, a similar approach is used in this study for setups that include a single flat plate.
 240 The mean drag coefficient of a cylinder with two parallel flat plates installed in the wake is
 241 modeled and compared in the second case, as shown in Figure 4.



242 FIG. 4) Numerical method validation for fixed cylinder and parallel plates. (Dehkordi and Jafari, 2010)

243 The flat plates are positioned at a vertical distance of $0.75H$ from the horizontal axis. Upon
244 reviewing the results, the reduction of the gap leads to a slow decrease in the $C_{d_{mean}}$. When
245 the gap reaches a magnitude of $G = 2.3$, the mean drag experiences a significant decrease,
246 followed by a slight reduction for smaller gaps. The little differences in results support the
247 accuracy of the numerical method in predicting the flow dynamics between entities and the
248 impact of downstream plates on the upstream cylinder. The numerical methodology used
249 demonstrates a high level of accuracy in simulating Flow-Induced Vibration (FIV) for either a
250 single downstream flat plate or two parallel plates.

251

252 3 Results and discussions

253 This section explains the details of Flow-Induced Vibration (FIV) on downstream single or
254 parallel flat plates. Initial simulations were conducted under different setups at a lower speed
255 of 6, where the natural frequency of the plate(s) downstream matched the vortex shedding
256 frequency of the cylinder upstream. The results are divided into three separate subsections
257 to enhance clarity and promote better comprehension. The first paragraph describes the
258 behavior of a single plate (Section 3.1), covering the vibration amplitude, applied forces, and
259 wake structure in various setups. It also outlines the response patterns for cases with high
260 amplitudes. The following part describes the results for the parallel plates in a consistent way
261 (part 3.2). A comparison analysis is conducted to determine the potential effectiveness of the
262 system in developing an energy harvesting mechanism (Section 3.3).

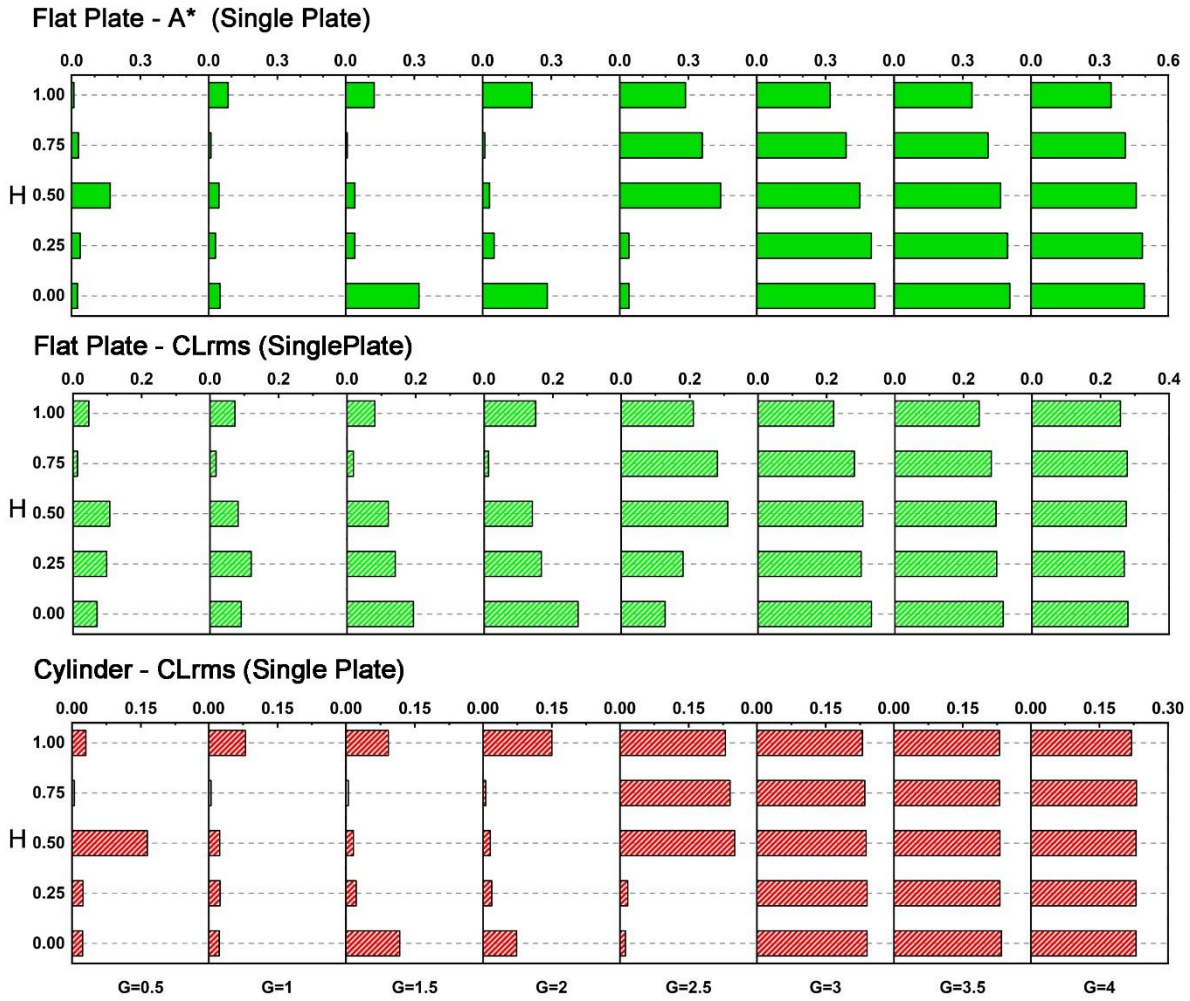
263 3.1 WIV of a single wake-mounted flat plate

264 This section delves into the study of Flow-Induced Vibration (FIV) of a single flat plate,
265 examining various configurations consisting of eight horizontal and five vertical locations. The
266 analysis involves a thorough evaluation of vibration amplitude and hydrodynamic forces, as
267 well as an investigation of wake structure using pressure coefficient, vorticity, and streamline
268 contours. The study then focuses more on cases with large amplitudes by including a range
269 of lower velocities.

270

271 3.1.1 Vibration amplitude and hydrodynamic forces

272 The vibration amplitude of the single flat plate in various horizontal and vertical positions is
273 illustrated in FIG. 5. As the lift force is usually the main source of energy for transverse
274 vibration, the root mean square values of lift coefficient (CL_{rms}) are also presented for both
275 objects. According to the results, the plate's response can be divided into two distinct parts
276 based on its amplitude, determined by a critical horizontal position of $G=3$.

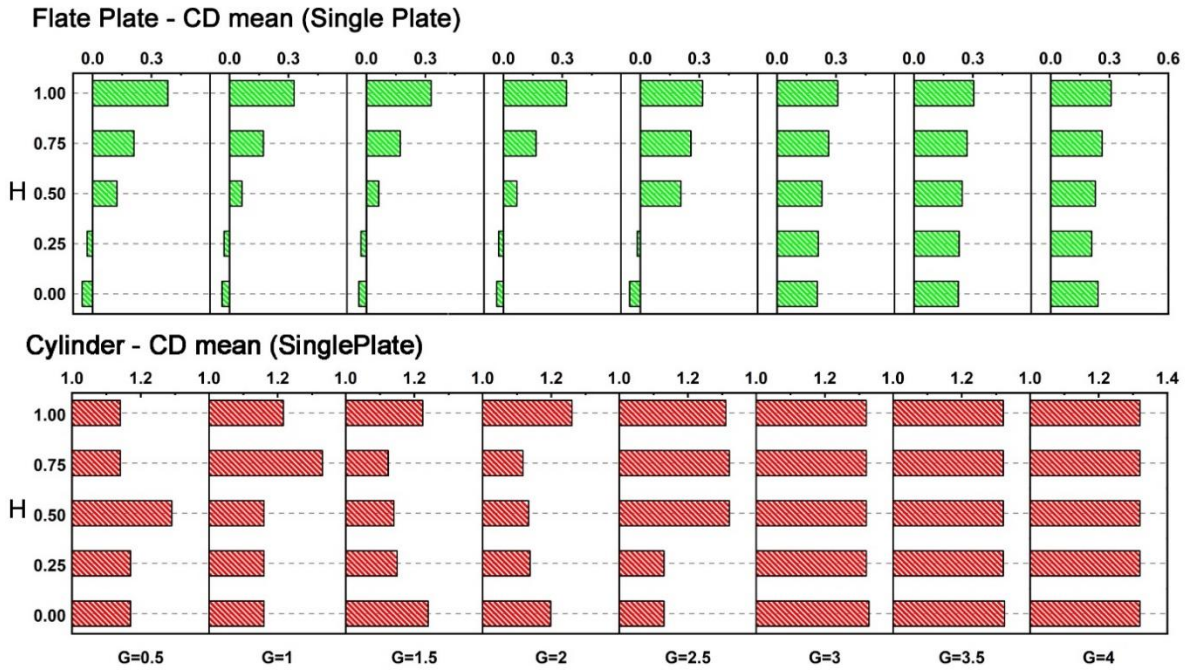


277 **FIG. 5) Vibration amplitude and CL_{rms} of the single flat plate and the cylinder in different configurations at $Ur=6$**
 278 When the plate is mounted at horizontal distances of $G \geq 3$, it causes oscillations at all vertical
 279 positions with a significant amplitude. The highest point of A^* occurs when the plate is in line
 280 with the center axis of the wake ($H=0$). When the vibration moves away from the central
 281 point, the amplitude decreases gradually until it reaches its lowest point at $H=1$. Conversely,
 282 when the horizontal spacing is less than 3 ($G < 3$), the plate shows a significant reliance on its
 283 placement. At $G=2.5$, there is minimal vibration at $H=0$. However, vertical separations greater
 284 than 0.25 cause a sudden increase in amplitude at $H=0.5$, followed by a gradual decrease. The
 285 amplitude is typically low and may be regarded insignificant for most designs across short
 286 horizontal intervals, except for $G=1.5, 2$. The horizontal positions are linked to moderate and
 287 somewhat reduced amplitudes at $H=0$ and 1, respectively. For the highest vertical separation

288 (H=1), the plate vibrates over a wider range of horizontal distances, reaching up to G=1.5. This
289 indicates the possible existence of a different vibratory mechanism, which will be further
290 investigated in the following section.

291 The variation of lift for the cylinder shows that placing the flat plate more than 3 times the
292 cylinder diameter away has no effect on the lift properties of the cylinder upstream,
293 regardless of its vertical position. A noticeable relationship is seen between the amplitude A^*
294 of the plate and the root-mean-square lift coefficient (CL_{rms}) of the cylinder when the
295 horizontal gap is limited. At G=2.5, there is a minimal increase in lift at H=0 and 0.25, but
296 substantially higher values are observed at higher vertical positions. The CL_{rms} of the cylinder
297 increases at narrower gaps, especially when the flat plate shows significant amplitude. A
298 significant coefficient of lift is observed for the flat plate when placed at $G \geq 3$. At these levels,
299 the lift slightly decreases as the plate is moved to higher vertical locations. Examination of
300 setups with a G factor of 2.5 shows that the highest lift is reached at a height of 2.5,
301 corresponding to the increase in vibrational magnitude. Instances with G values of 1.5 and 2
302 exhibit similar patterns, showing increased lift at H=0 that decreases at higher vertical
303 locations and becomes insignificant at H=0.75, before increasing again at H=1. With smaller
304 horizontal distances, the CL_{rms} often stays relatively low.

305 The impact of the moving flat plate on the flow patterns behind the cylinder is also evident in
306 the average drag force. For $G \geq 3$, the plate has minimal effect on the cylinder, as seen in FIG.
307 6. An increase in the mean drag coefficient is noticed for the cylinder at G=2.5-H=0.5, however
308 at closer horizontal positions, the flat plate mostly causes a decrease in the drag force on the
309 cylinder.



310 FIG. 6) Mean values of drag coefficient for the cylinder and single flat plate in different configurations at $Ur=6$

311 The drag coefficient mean of the flat plate increases as it moves away from the wake
 312 centerline at all horizontal intervals. While the increase in rate is considered low for $G \geq 3$, it
 313 becomes more noticeable in smaller gaps, showing a decrease in average drag at $H=0, 0.25$.
 314 This suggests possible changes in flow patterns with adjustments in vertical placement.

315 Examining the lift and drag coefficients confirms the previous claim that a horizontal
 316 separation of $G=3$ is a crucial boundary. Oscillation occurs inside this boundary, regardless of
 317 vertical alignment, and has minimal impact on the lift and drag forces of the upstream
 318 cylinder.

319 3.1.2 Wake Structure and vibration mechanism

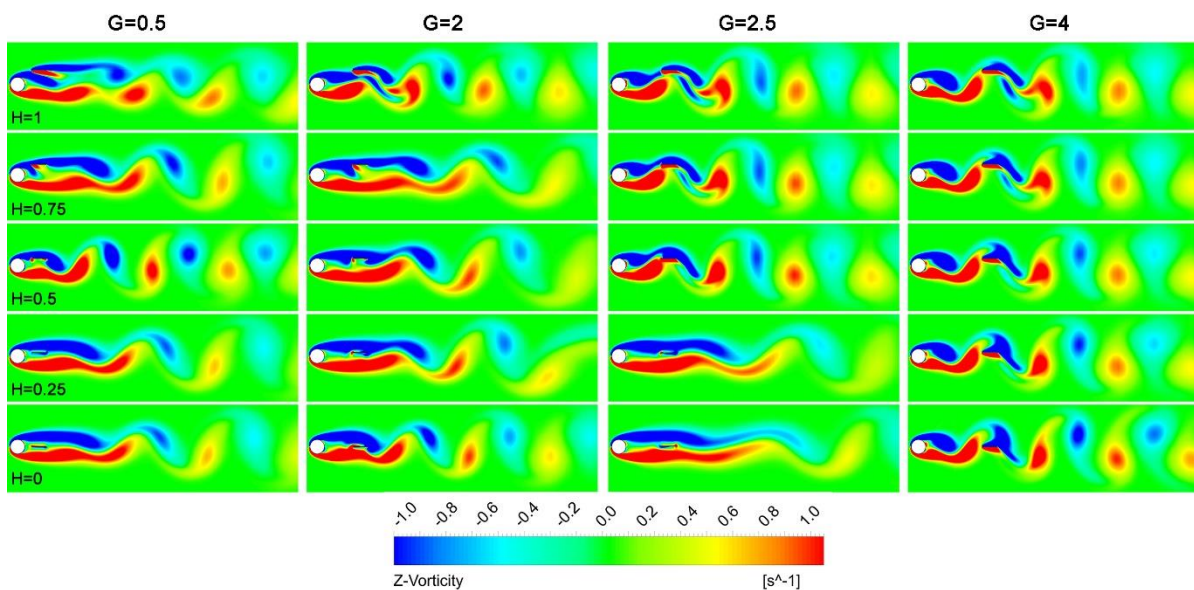
320 The system's response can be divided into four separate parts depending on the plate's
 321 amplitude, as outlined in TABLE IV. The first segment includes situations where the plate is
 322 attached at $G=0.5, 1$, showing modest or insignificant amplitudes. The following section
 323 includes cases with G values of 1.5 and 2, where significantly high amplitudes are recorded

324 when the plate is located at heights of 0 and 1. At $G=2.5$, there is a noticeable increase in A^*
 325 at $H=0.5$. The section from $G=3-4$ demonstrates consistently high A^* values at all elevations.

326 **TABLE IV) The system response of a single plate categorized by vibration amplitude.**

Horizontal Distance	Main Characteristics
$G= 0.5, 1$	Small A^* regardless of H
$G= 1.5, 2$	Relatively large A^* at $H=0, 1$
$G= 2.5$	Sudden jump in A^* at $H=0.5$
$G= 3, 3.5, 4$	Large A^* regardless of H

327 Instantaneous non-dimensional vorticity contours are illustrated at the moment when the
 328 plate reaches the highest point of its vibratory cycle to thoroughly analyze the flow dynamics
 329 in each section. FIG. 7 shows the wake arrangement for G values of 0.5, 2, 2.5, and 4,
 330 representing examples for each segment.



331 **FIG. 7) The wake structure for all configurations with $G=0.5, 2, 2.5, 4$**

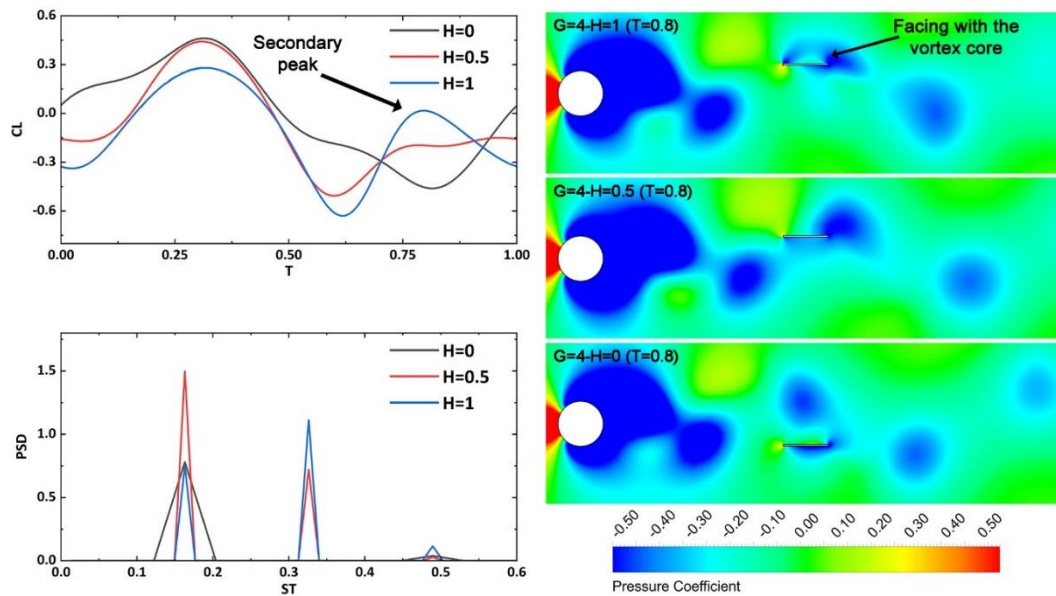
332 In the case with $G=0.5$, a consistent pattern is shown for $H=0$ and 0.25 . The flat plate causes
 333 the creation of elongated and distinct shear layers, similar to a stationary splitter. This
 334 configuration delays the start of vortex shedding to a location further downstream, resulting
 335 in a lower frequency and confirming the minimum values of CL_{rms} for the objects in question.

336 This wake arrangement provides additional evidence for the small values of CD_{mean} . When
337 the flat plate is raised to a higher vertical position at $G=0.5-H=0.5$, vortex shedding begins at
338 the trailing edge of the plate. The vortex cores are closer together on a horizontal level,
339 causing a greater change in lift. This results in a higher frequency and, as a result, a larger
340 vibration amplitude for the plate. When the flat plate is moved to $H=0.75$, a different wake
341 structure is revealed, with the upper shear layer splitting into two segments. The proximal
342 portion interacts with the lower shear layer, while the distal segment of the upper layer, along
343 with the lower one, delays vortex shedding further downstream, similar to the conditions
344 seen in the $H=0, 0.25$ cases. At the maximum vertical distance ($H=1$), the splitting of the upper
345 shear layer is clearly visible, with a larger proximal section. In these last two setups, the flat
346 plate is mainly located inside or even beyond the boundaries of the upper shear layer, facing
347 a flow with increased velocity and energy, resulting in a higher mean drag coefficient.

348 When the flat plate is positioned along the central axis of the wake ($H=0$) in the second case
349 ($G=2$), the shear layers converge near the end of the plate, leading to the initiation of vortex
350 shedding. This design is similar to the one seen at $G=0.5-H=0.5$. The following interaction
351 increases lift fluctuations at a frequency close to the natural frequency, resulting in a rise in
352 vibration amplitude. Further explanation of this phenomenon may be found in the author's
353 previous study (Jebelli and Masdari, 2022). Increasing the vertical spacing to $H=0.25$ or 0.5
354 causes elongated shear layers to reappear, with the flat plate acting as a splitter. The upper
355 shear layer bifurcates when $H=0.75$, but at $H=1$, a different flow dynamic occurs. A larger
356 portion of the upper shear layer consistently divides, leading to an earlier start of vortex
357 shedding compared to $G=0.5-H=1$ in the nearby wake. This wake setup increases lift and drag
358 coefficients, while also improving the vibratory sensitivity of the plate.

359 The flow dynamics at $G=2.5$ require significant attention, especially because of the critical
360 point at $G=2.5-H=0.5$. When $H=0.25$, two elongated shear layers reappear, and vortex
361 shedding is delayed with a significantly lower frequency. Placing the flat plate at the crucial
362 point of $H=0.5$ results in a significant change in the wake pattern, leading to the initiation of
363 vortex shedding between the two objects. This modification enhances the periodic flow
364 setup, increasing the CL_{rms} and average drag on both objects, and amplifying the plate's
365 vibrational reaction. Moving the flat plate to higher vertical positions ($H=0.75, 1$) results in a
366 matching wake formation. However, when the distance increases between the plate's
367 location and the central axis of the wake, the plate's interaction mostly focuses on the vortices
368 coming from the top shear layer.

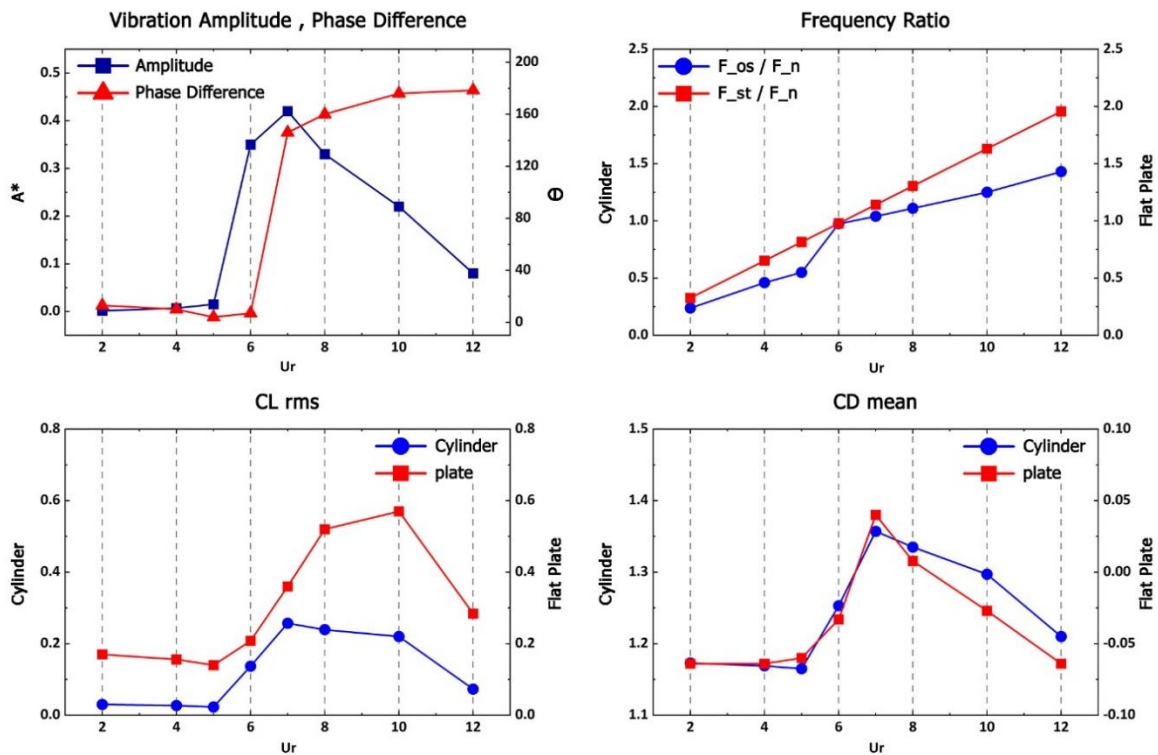
369 Placing the plate with a significant horizontal distance ($G=4$) creates enough space between
370 the objects, which helps the shear layers interact and leads to the shedding of vortices at
371 various vertical positions. In such configurations, the impact of the plate on the cylinder
372 located upstream is negligible, resulting in nearly identical lift and drag forces. An increase in
373 vertical separation results in a gradual rise in drag and a corresponding decline in lift forces
374 for the flat plate. The impact of the upper vortex becomes more noticeable as the altitude of
375 H increases. Figure 8 displays pressure coefficient contours, lift coefficient fluctuations, and
376 dominant frequencies for a single vibratory cycle at $G=4-H=0, 0.5, 1$.



377 **FIG. 8)** Variation of lift coefficient, dominant frequencies and pressure coefficient contours for $G=4-H=0, 0.5, 1$
 378 At $H=0$, a unique frequency corresponding to the vortex shedding frequency of the cylinder is
 379 observed. As the plate's vertical position increases, a secondary frequency, twice the initial
 380 one, appears and grows stronger. At $H=0.5$, the primary frequency stays the same, but when
 381 the plate is mounted at $H=1$, the secondary frequency takes over. This phenomenon is also
 382 seen in the lift variation, where a secondary peak appears during one vibratory cycle. Studying
 383 the pressure coefficient contours at $T=0.8$ reveals the differences and origin of the secondary
 384 frequency.
 385 During one complete vibration cycle at $H=0$, the plate does not interact with the vortex cores
 386 directly. Each vortex passes by while the plate is at the furthest point of its vibrating range.
 387 By increasing the vertical spacing to $H=1$, the shedding vortex from the higher shear layer
 388 meets the plate at a point where the central mounting and vibrational amplitude prevent the
 389 plate from moving down enough to avoid a direct collision with the vortex. As a result, a low-
 390 pressure area forms above the plate, leading to a higher lift coefficient compared to what is
 391 seen at ground level ($H=0$).

392 **3.1.3 Response Branch**

393 The main goal of this study is to analyze various configurations related to Wake-Induced
 394 Vibration (WIV) of flat plates. The paper conducts a detailed analysis of configurations that
 395 show significant vibrational amplitudes. The instance where $G=1.5$ and $H=0$, showing a large
 396 amplitude (A^*) with a small horizontal separation, is described as the first case. The scenario
 397 with the greatest separation and significant amplitude, while not affecting the upstream
 398 cylinder ($G=4-H=0$), is chosen for further analysis. Figure 9 explains the results related to
 399 $G=1.5-H=0$, which have been expanded to include a range of lower speeds.



400 **FIG. 9) WIV of the single flat plate at $G=1.5-H=0$ in a range of reduced velocities**

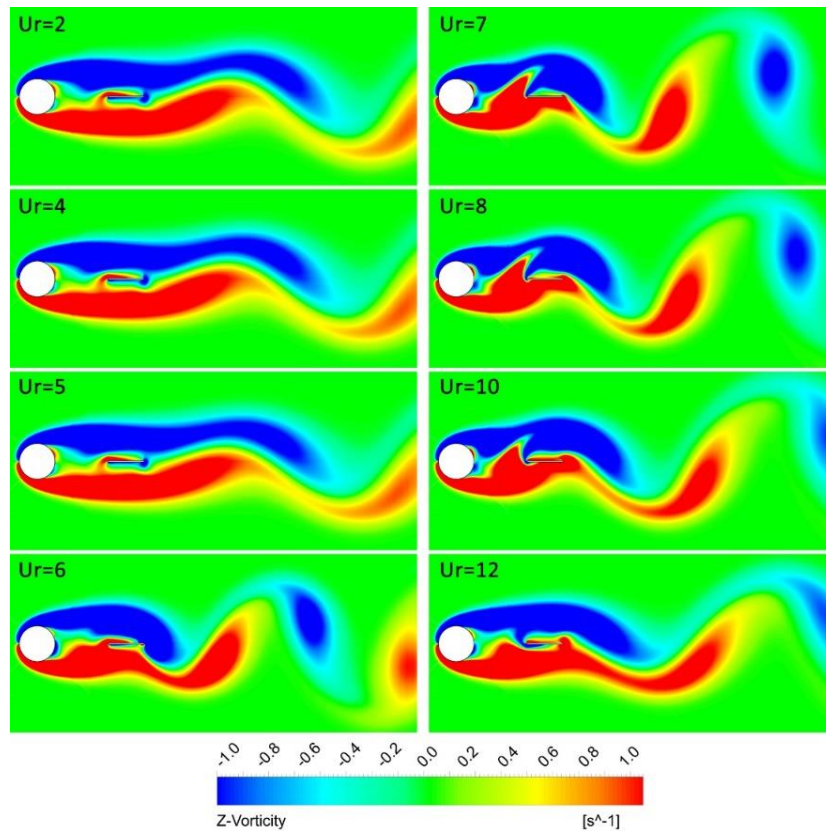
401 The experiment shows that the vibration amplitude stays insignificant throughout the
 402 restricted velocity range of $Ur=2-5$. As the lowered velocity increases, there is a significant
 403 increase in amplitude, reaching its peak at $Ur=7$. The vibratory response decreases linearly till
 404 reaching $Ur=12$. A^* at $Ur=12$ significantly surpasses those observed at $Ur=2-5$. The phase
 405 disparity study shows that the plate's vibration and lift force are synchronized for $Ur=2-6$, with

406 a sudden change in Θ occurring at $Ur=7$. At high speeds, the angle Θ steadily increases,
407 approaching 180 degrees. This phase transition is similar to what is seen in Vortex-Induced
408 Vibrations (VIV) of circular cylinders, happening when the vibratory frequency exceeds the
409 natural frequency ($F/F_n > 1$).

410 An analysis of frequency ratios shows a gradual increase in vibration within the velocity range
411 of 2-5, although at a slower rate compared to the vortex shedding frequency of a cylinder
412 without any obstacles. At $Ur=6$, the vibratory frequency experiences a notable increase,
413 matching the Strouhal number (F_{st}). This rise in height happens simultaneously with the
414 increase in vibration amplitude of the plate. At higher decreased velocities, there is a linear
415 increase in frequency ratio, although the rate of growth is slower than that of the F_{st} . The
416 plate always reflects the vortex shedding frequency, and its presence alters both the wake
417 structure and the vortex shedding frequency of the cylinder ahead. Both the cylinder and the
418 plate show a similar trend in the lift coefficient, with an increase starting at $Ur=6$, stable values
419 between $Ur=8-10$, and a reduction at $Ur=12$. The average drag coefficient also shows a pattern
420 of both rise and reduction. A negative mean drag coefficient is observed for the flat plate for
421 most decreased velocities, except for $Ur=7$ and 8.

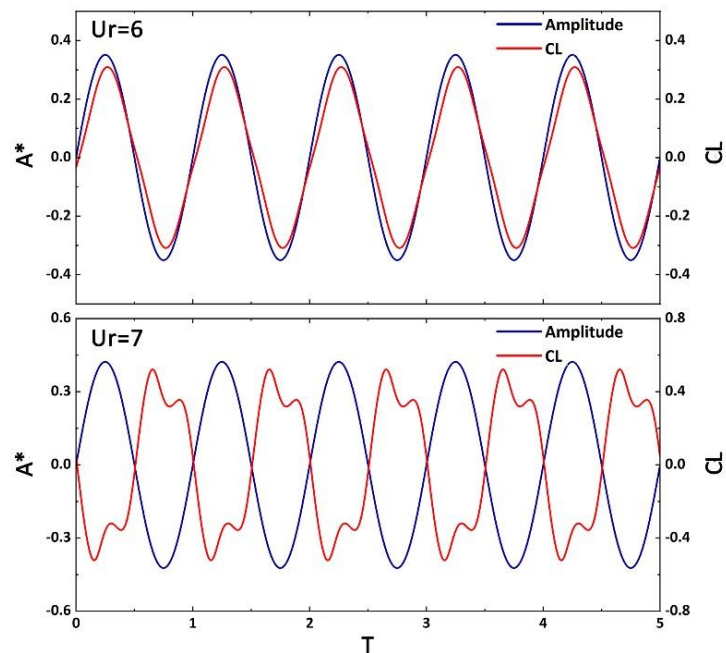
422 Figure 10 shows the wake creation for this particular configuration at different decreased
423 velocities when A^* equals 0. The immediate vorticity patterns reveal a consistent wake
424 structure for flow velocities ranging from 2 to 5, corresponding with the vibration and lift
425 characteristics. Adding the plate delays the start of vortex shedding to a location further down
426 the stream, creating a balanced flow pattern around the objects, resulting in low CL_{rms} values.
427 This setup leads to a decrease in the average drag coefficient, with negative coefficients seen
428 for both the cylinder and the plate. Vortex shedding begins at $Ur=6$ at the end of the flat plate,

429 when the plate's inherent frequency closely aligns with the shedding frequency. Furthermore,
 430 certain interactions can be observed at this speed between the shear layers in the gap.



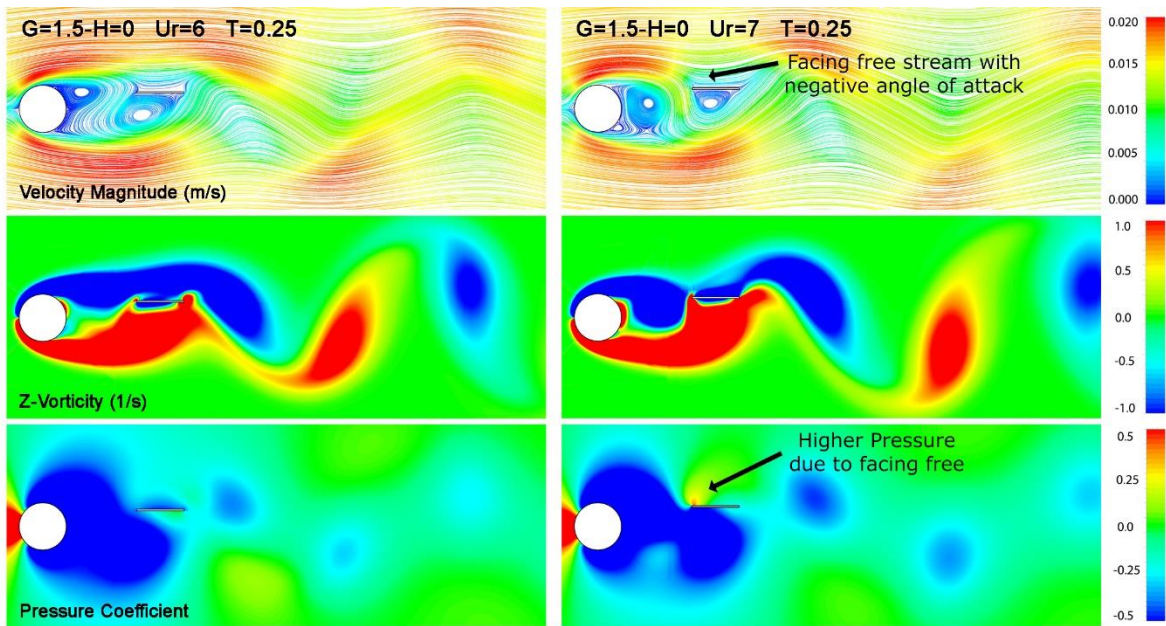
431 **FIG. 10) Instantaneous vorticity contours in different reduced velocities for $G=1.5-H=0$**
 432 At $Ur=7$, the structure of the near-wake undergoes changes, with increased interaction
 433 between shear layers and a variation in the vortex production pattern. As the plate moves up
 434 or down, the shear layer forms nearer to the back of the cylinder. This setup helps the vortex
 435 detach from the shear layer on one side of the plate as it evolves and surrounds the other
 436 side. The maximum mean drag coefficient is observed at this velocity due to the low-pressure
 437 shear layer at the rear of the cylinder, as shown in FIG. 9. At $Ur=8$ and 9, the vortex production
 438 process continues but with reduced intensity. However, at $Ur=10$, the interaction in the wake
 439 decreases significantly, and the wake formation returns to features seen at lower reduced
 440 velocities.

441 FIG. 11 illustrates the temporal fluctuation of vibration amplitude and lift coefficient at $Ur=6$
 442 and 7 to explain the flow dynamics for velocities with increasing amplitudes and phase shift.
 443 The vorticity, streamlines, and pressure coefficient contours are shown at $T=0.25$ of a
 444 vibratory cycle when the plate is at its maximum position, as illustrated in FIG. 12.



445 FIG. 11) Time variation of vibration amplitude and lift coefficient of the single flat plate for $G=1.5-H=0$

446 Vortices are observed to form in the near-wake structure at $Ur=7$, originating at the back of
 447 the cylinder. The plate's vertical movement causes the shear layer to start and develop
 448 toward the back of the cylinder, while the shear layer on the other side widens and covers
 449 the lower/upper part of the plate.



450 FIG. 12) The vorticity, streamlines and pressure coefficient contours for $G=1.5-H=0$ at $Ur=6, 7$ ($T=0.25$)

451 As the plate moves to the other side, it crosses the shear layer, causing the vortex to detach
 452 and go downstream. The plate moves beyond the shear layers and is exposed to the external
 453 flow at its highest points during this phase. After passing over the cylinder, the external flow
 454 hits the plate at an angle, creating a high-pressure area near the front edge, resulting in a
 455 downward force on the plate. This situation creates a significant difference in timing between
 456 the applied force and the resulting vibration.

457

458 The vibration amplitude increases significantly as the horizontal distance extends across all
 459 vertical alignments. Figure 13 illustrates the plate's response with $G=4$ and $H=0$.

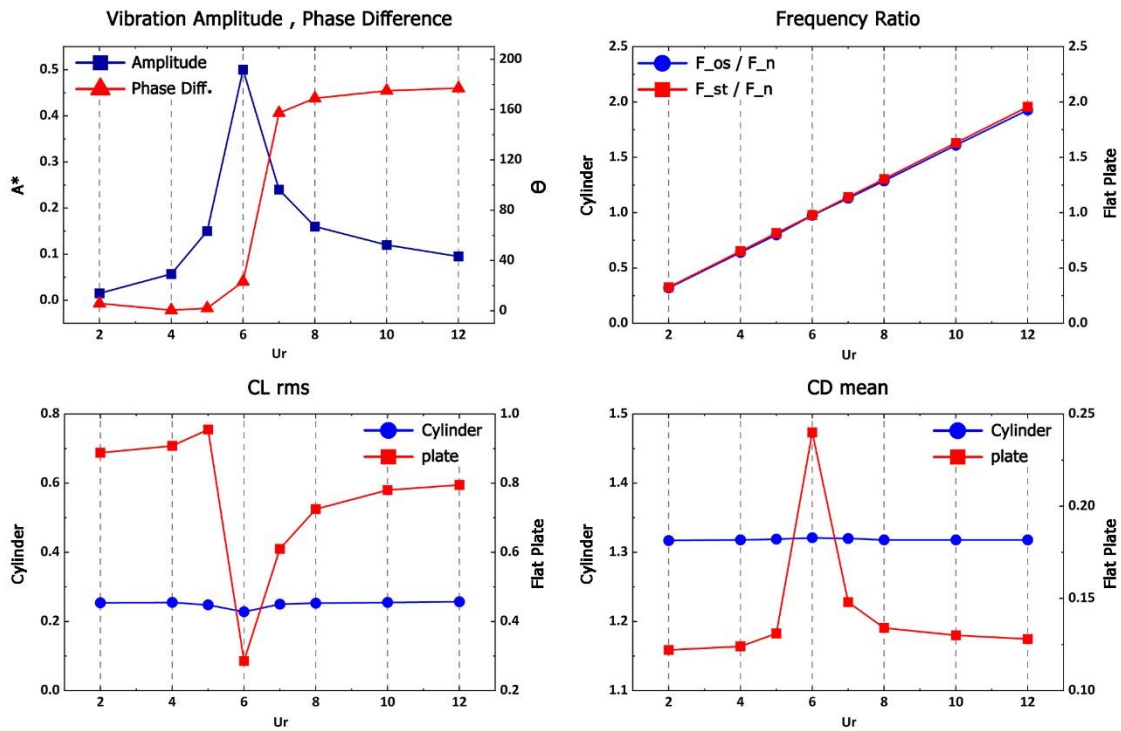


FIG. 13) WIV of the single flat plate at $G=4-H=0$ in a range of reduced velocities

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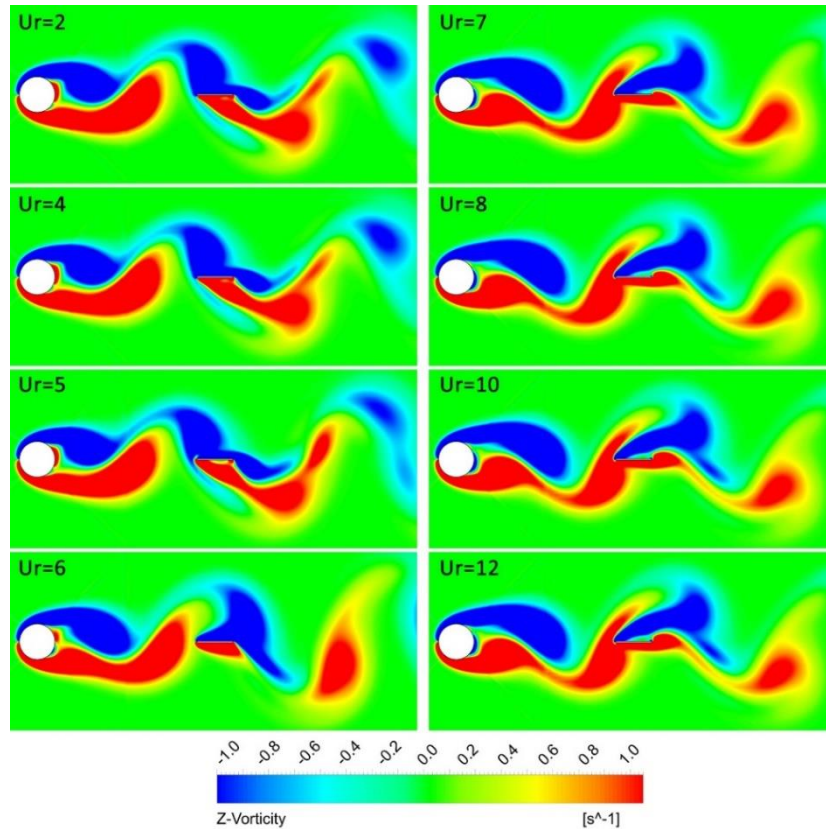
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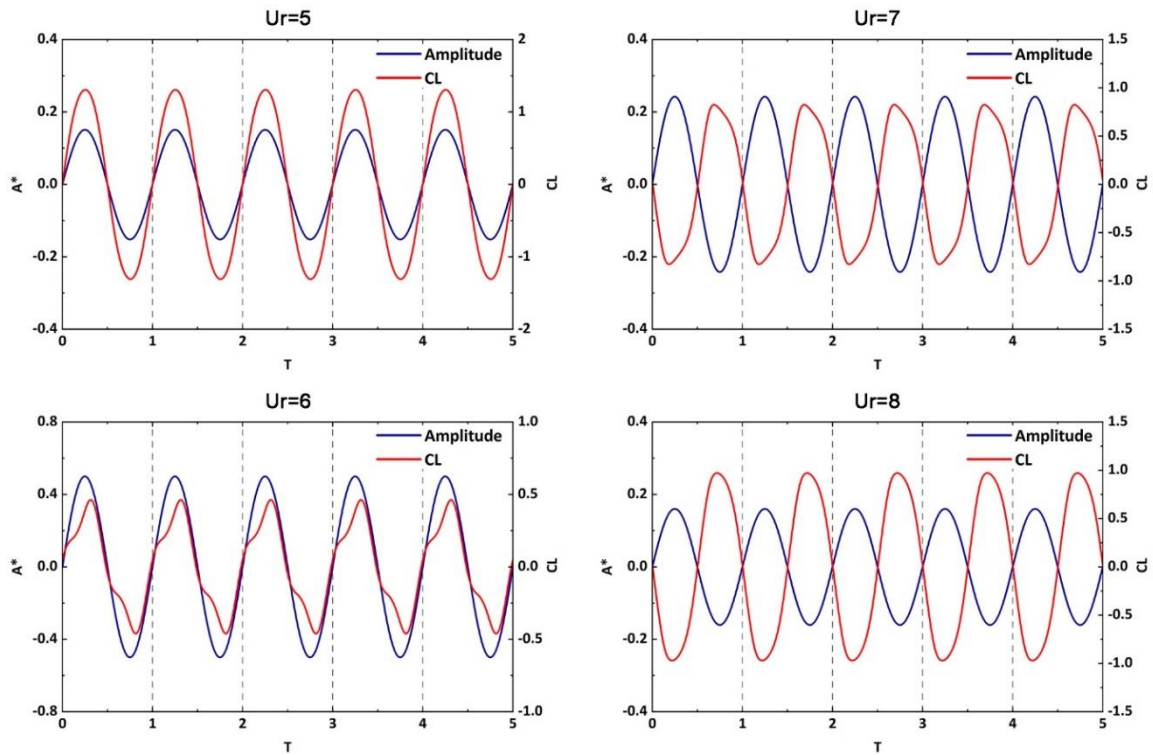
The amplitude of plate vibration increases gradually as the decreased velocities range from $U_r=2$ to 5. At $U_r=6$, there is a significant increase to a peak amplitude value of $A^*=0.5$, which is the highest among the other single plate setups. The amplitude decreases considerably at first, then stabilizes at greater velocities. The phase difference between force and vibration is similar to that seen in the $G=1.5-H=0$ setup, showing synchrony for U_r values between 2 and 5, a phase shift at $U_r=7$, and phase discordance at higher speeds. A sharp change in phase difference occurs when F/F_n is greater than 1, similar to what happened before. The analysis of frequency ratios shows that the plate consistently matches the vortex shedding frequency of the upstream cylinder at all velocities examined. Furthermore, it is clear that the plate's position far from the cylinder does not affect the vortex shedding frequency, unlike in earlier cases. The significant increase in amplitude at $U_r=6$ results in a noticeable drop in lift for the plate, as shown by the abrupt decline and subsequent increase in the CL_{rms} and mean drag of

473 the plate, respectively. The fluctuations in force confirm that the plate has little effect on the
474 upstream cylinder, with uniform values at all lower velocities.

475 Positioning the plate at $G=4-H=0$ allows sufficient space for the shear layers to interact, as
476 shown in FIG. 14. A consistent pattern of vortex shedding forms behind the cylinder as a
477 result.

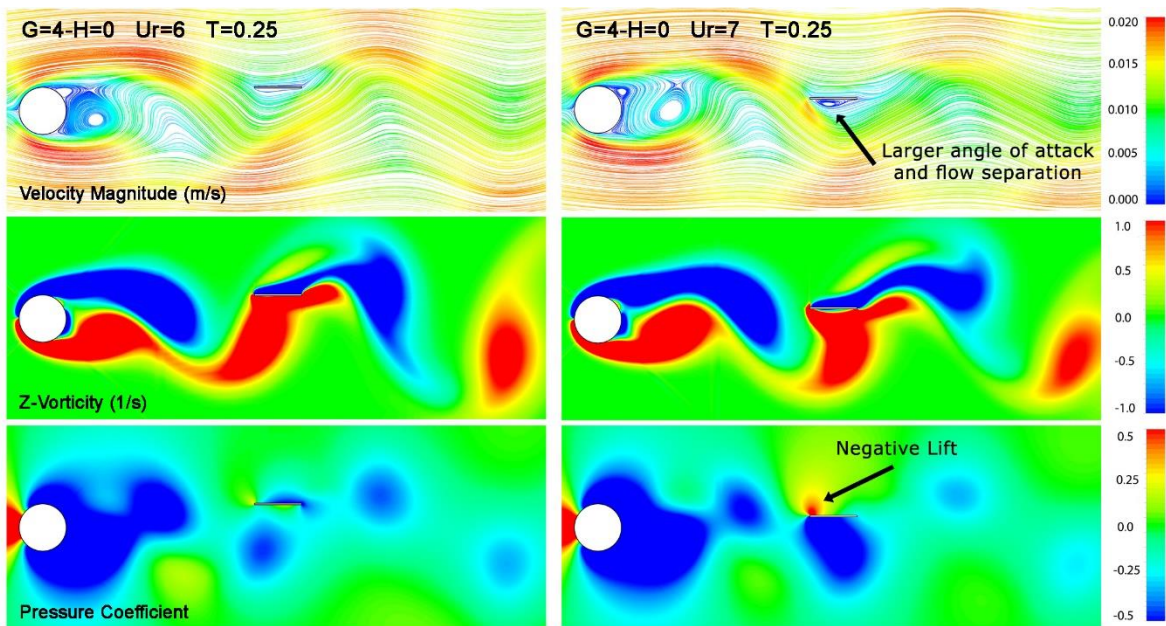


478 **FIG. 14)** Instantaneous vorticity contours in different reduced velocities for $G=4-H=0$
479 The plate interacts with oncoming vortices, which apply periodic forces, leading to the
480 increased lift forces observed. This wake structure leads to a significant increase in the drag
481 coefficient of the plate because it is no longer positioned between the shear layers and is
482 directly exposed to the vortices. At $Ur=6$, there is a significant increase in amplitude and a
483 decrease in CL_{rms} , followed by a phase transition at $Ur=7$. FIG. 15 illustrates the chronological
484 patterns of vibration and lift coefficient for $Ur=5-8$, aiding in a more profound understanding
485 of these occurrences.



486 **FIG. 15) Variation of vibration amplitude and lift coefficient of the single flat plate for $G=4-H=0$ ($Ur=5-8$)**

487 The lift coefficient's behavior depends on the lowered velocity, and there is a noticeable
 488 change in the CL profile as the Ur increases. The lift coefficient displays periodic behavior in
 489 all cases, but deviates from a sinusoidal pattern at lower velocities of $Ur=6$ and $Ur=7$. This
 490 departure is most noticeable at $Ur=6$. An increase in Ur results in a change in the size of the
 491 phase difference. At $Ur=7$ and 8 , the phase difference approaches around 180 degrees, while
 492 at $Ur=5$, it closely approximates zero. Figure 16 illustrates the wake configuration at $T=0.25$
 493 to highlight the differences at $Ur=6$ and 7 that cause the phase shift.



494 FIG. 16) The vorticity, streamlines and pressure coefficient contours for $G=4-H=0$ at $Ur=6, 7$ ($T=0.25$)

495 Although the wake structure is generally similar, especially in terms of vorticity, there are
 496 subtle differences that explain the phase shift mentioned earlier. As shown in FIG. 13, the
 497 vibration at various decreased velocities corresponds to the vortex shedding frequency
 498 produced by the cylinder. Therefore, as the speed decreases, the lift force frequency deviates
 499 from the natural frequency, resulting in a decrease in vibration amplitude at $Ur=7$. The
 500 decreased maximum amplitude causes the plate to experience a higher angle of attack when
 501 it encounters the external flow, resulting in higher velocity and energy, which leads to flow
 502 separation on the lower side of the plate. This setup creates an increased pressure difference
 503 and a stronger vertical force, directed opposite to the motion. This process explains the phase
 504 difference seen at $Ur=7$.

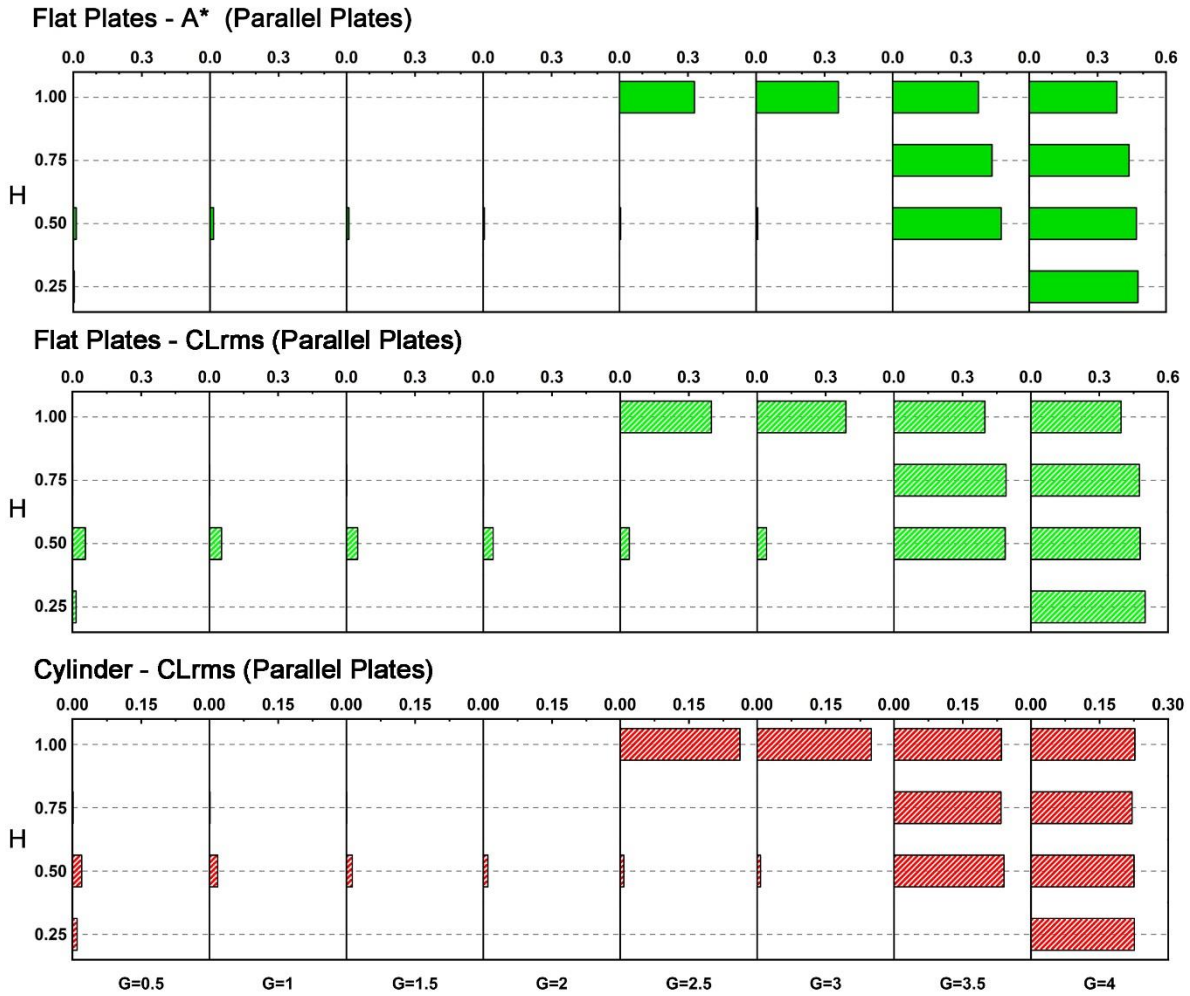
505 Due to the proven effectiveness of a single plate in Wake-Induced Vibrations (WIV) regardless
 506 of its position on the wake centerline, parallel plates have been selected for further
 507 investigation. More information on this will be provided in the next section.

508 3.2 WIV of parallel wake-mounted flat plates

509 This part examines the Wake-Induced Vibrations (WIV) of two parallel flat plates, analyzing
510 eight horizontal and four vertical configurations. The examination includes an analysis of the
511 plates' vibration amplitude and hydrodynamic forces, with the results being properly
512 reported. The research also investigates wake dynamics by analyzing pressure coefficient,
513 vorticity, and streamline contours. The study further examines the plates' vibratory behavior
514 and potential reaction branches in scenarios with substantial vibration amplitude, following
515 the same approach as the previous section.

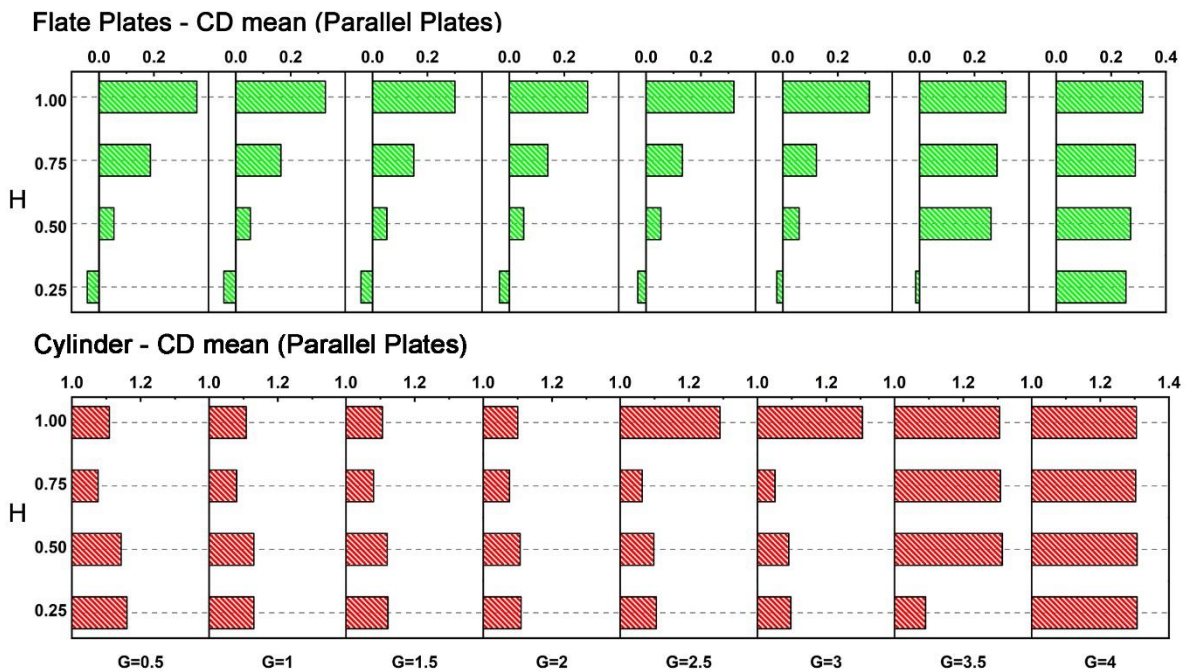
516 3.2.1 Vibration amplitude and hydrodynamic forces

517 FIG. 17 shows the vibration amplitude of the flat plates installed in parallel with the wake, as
518 well as the change in force in the transverse direction, for all objects. The amplitude measured
519 on the plates indicates low vibration in the nearby airflow at all vertical levels. However,
520 significant vibrations are detected when the plates are positioned far apart both horizontally
521 and vertically. Vertical alignments within the range of $G=0.5-2$ exhibit insignificant
522 amplitudes. The zone without vibration extends from $G=2.5$ to 3 , except for $H=1$, where
523 exceptionally large amplitudes are recorded. At $G=3.5$, the plates start vibrating for vertical
524 separations of $H \geq 0.5$, and choosing the maximum horizontal separation ($G=4$) causes
525 vibration in all setups. An increase in the vertical distance between the plates leads to a
526 gradual decrease in the peak amplitude, a trend also observed in setups with a single flat
527 plate. Another important feature is the appearance of a smaller vibration-free zone that
528 occurs at greater vertical distances, which is also noticeable in the situation with just one
529 plate.



530 **FIG. 17) Vibration amplitude and CL_{rms} of the parallel flat plates and the cylinder in different configurations at $Ur=6$**
 531 The CL_{rms} of the upstream cylinder shows correlations with the amplitude of the plates,
 532 making it easier to identify areas that either show vibration or do not. In areas where the
 533 plates are vibrating, the lift force on the cylinder is high and consistent, indicating that the
 534 plates have no effect on it. When the plates are in equilibrium, the CL_{rms} significantly
 535 decreases, indicating a form of wake flow modulation. The CL_{rms} of the plates also follows this
 536 pattern, showing a gradual reduction in lift as the vertical distance increases. A little increase
 537 in the lift coefficient was seen for objects located in the non-vibratory zone, particularly at a
 538 vertical position of $H=0.5$.

539 The variations in the mean drag coefficient show similarities to the vibration amplitude,
 540 especially with the cylinder, but noticeable differences are also apparent, especially with the
 541 flat plates (FIG. 18). The cylinder shows a consistently even CD_{mean} when the plates are placed
 542 at appropriate distances, both horizontally and vertically. An increased average drag is
 543 noticed when the plates are located outside the non-vibratory area. On the other hand,
 544 significant decreases in CD_{mean} are observed in different setups.



545 **FIG. 18)** Mean values of drag coefficient for the cylinder and parallel flat plates in different configurations
 546 Flat plates have a negative average drag when positioned at the lowest vertical point ($H=0.25$)
 547 for most horizontal separations, except for $G=4$, where a higher positive value is observed.
 548 The CD_{mean} for the parallel plates mirrors the behavior seen with a single plate. When the
 549 plate(s) are located on or near the wake centerline without vibrating, a decrease in average
 550 drag is seen, and raising the plate vertically results in greater values. The CD_{mean} of the plates
 551 shows a progressive increase at $G=4$, unlike the sharp change from negative to significantly
 552 higher values observed in other setups ($G=0.5-3.5$).

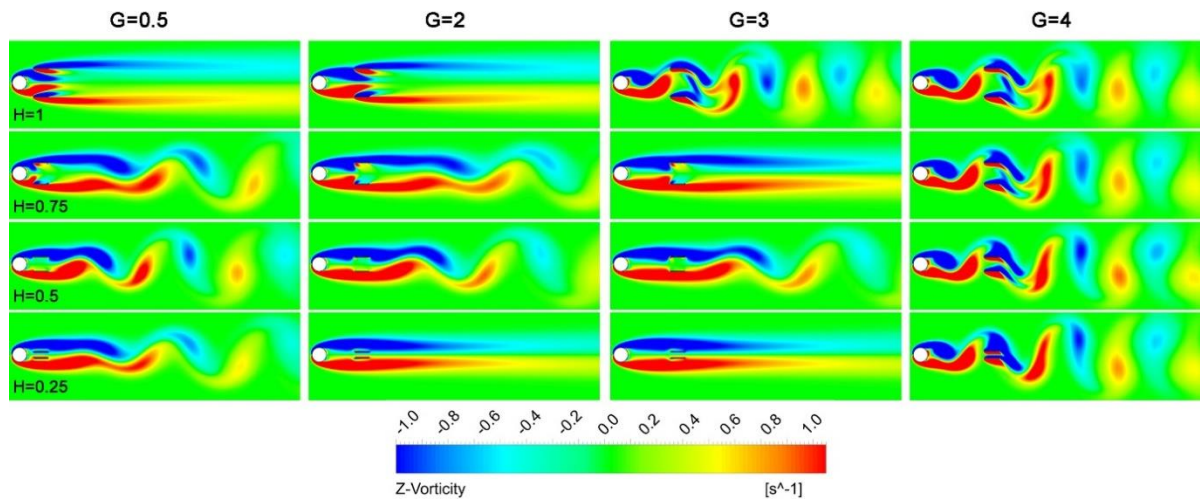
553 **3.2.2 Wake structure**

554 Based on the findings of the previous section, the amplitude response of the parallel plates
 555 can be divided into three specific groups (TABLE V). In the first group ($G=0.5-2$), there is no
 556 significant vibration measured at any vertical point, indicating a controlled wake flow. In the
 557 second group ($G=2.5-3$), vibration starts only at $H=1$. Plates in the final group ($G=3.5, 4$) display
 558 vibratory behavior regardless of the vertical separation, except for $G=3.5-H=0.25$.

559 **TABLE V) The system response of parallel plates categorized by vibration amplitude.**

Horizontal Distance	Main Characteristics
G= 0.5, 1, 1.5, 2	Negligible A^* regardless of H
G= 2.5, 3	Large A^* only at $H=0, 1$
G= 3.5, 4	Large A^* regardless of H

560 The non-dimensional vorticity contours are displayed to analyze the flow dynamics at the
 561 moment when the plates reach their vibrational midpoint. FIG. 19 shows the wake formation
 562 for G values of 0.5, 2, 3, and 4, indicating different stages of the process. Configurations with
 563 G values between 0.5 and 2 show low amplitude, and the plates remain stable. The results
 564 show that the addition of the plates causes different flow control methods, depending on the
 565 vertical distance between them. At lower intervals of $H=0.25$ and 0.5 , the plates completely
 566 separate the shear layers. Choosing $H=0.75$ leads to the formation of bifurcated layers. The
 567 interior segments are concentrated in the proximal wake, while the rest are isolated and start
 568 interacting further downstream. At the highest vertical distance ($H=1$), bifurcated shear layers
 569 reappear, with the inner sections more prominent and the remaining parts creating two
 570 elongated, separate layers without contact.



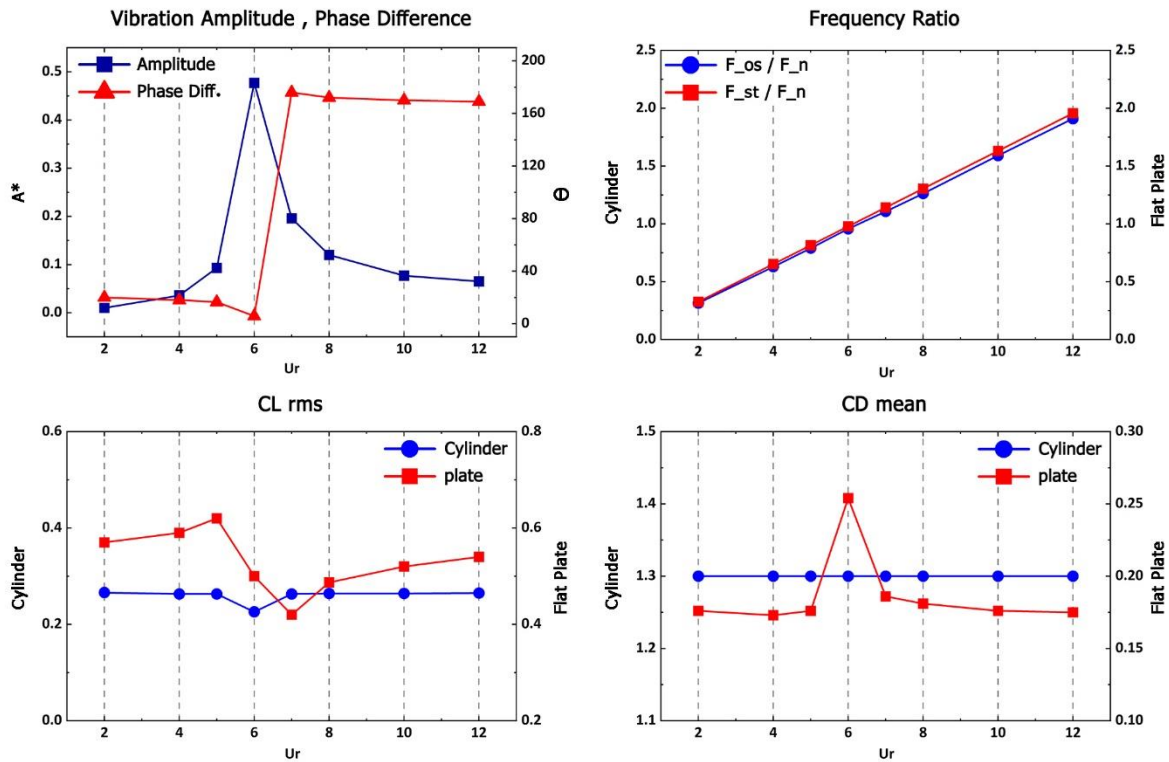
571 FIG. 19) The wake structure around the cylinder and parallel plates for all configurations with $G=0.5, 2, 3, 4$
 572 Within the intermediate range ($G=2.5, 3$), the plates show stability with negligible vibration
 573 at three vertical positions ($H=0.25, 0.5, 0.75$), and the wake configuration closely resembles
 574 that seen in the next segment. Increasing the vertical distance to $H=1$ allows enough space
 575 for the shear layers to come together and merge. Therefore, the plates experience shedding
 576 vortices, creating an uneven flow field. Choosing the maximum horizontal separation ($G=4$)
 577 causes vibration in all vertical alignments. Vortex shedding begins in the proximal wake and
 578 in front of the parallel plates, creating an uneven flow field and causing a significant variation
 579 in the lift coefficient for all objects.

580 At G values between 0.5 and 3 and H value of 0.5, the CL_{rms} is slightly higher for both entities
 581 in the vibration-free zone. An analysis of the vorticity patterns in this span reveals that vortex
 582 shedding begins quite close, even though it starts later downstream in all configurations. The
 583 initial interaction with the parallel plates causes changes in the formation of shear layers,
 584 leading to more significant oscillations in the lift force.

585 3.2.3 Response Branch

586 Upon examining the amplitude response of the plates, it is noted that A^* is significant
 587 regardless of the vertical separation at $G=4$. Therefore, two configurations, $G=4-H=0.25$ and

588 G=4-H=1, have been selected for further detailed examination, with simulations extended to
 589 include a range of lower velocities. The second part showed similarities in the development
 590 of waves in both of these environments. Therefore, the results and relevant debates will be
 591 presented simultaneously. FIG. 20 shows the results for the smaller vertical distance (H=0.25).

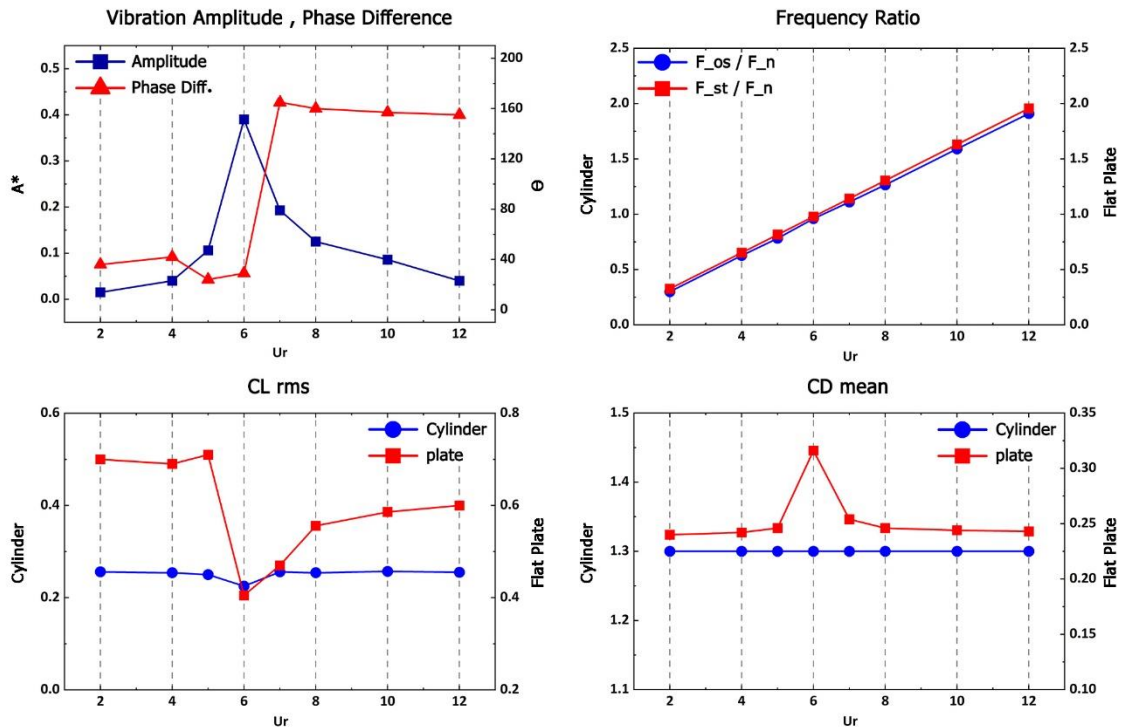


592 FIG. 20) FIV of the parallel flat plates for G=4-H=0.25 in a range of reduced velocities

593 The vibration amplitude gradually rises at progressively decreased velocities between Ur=2-
 594 5, then sharply increases to a peak value of $A^*=0.48$ at Ur=6, which is the highest amplitude
 595 recorded in parallel plates. At higher decreased velocities, there is an initial decrease in
 596 amplitude, followed by a little reduction in A^* .

597 At lower velocities between Ur=2 and 5, the lift force and vibration are coordinated. At Ur=7,
 598 a phase transition happens where vibrations go out of phase at high speeds. This phase
 599 change occurs repeatedly when the frequency ratio surpasses 1. The analysis of the frequency
 600 ratios shows that the plates reliably match the vortex shedding frequency of the upstream

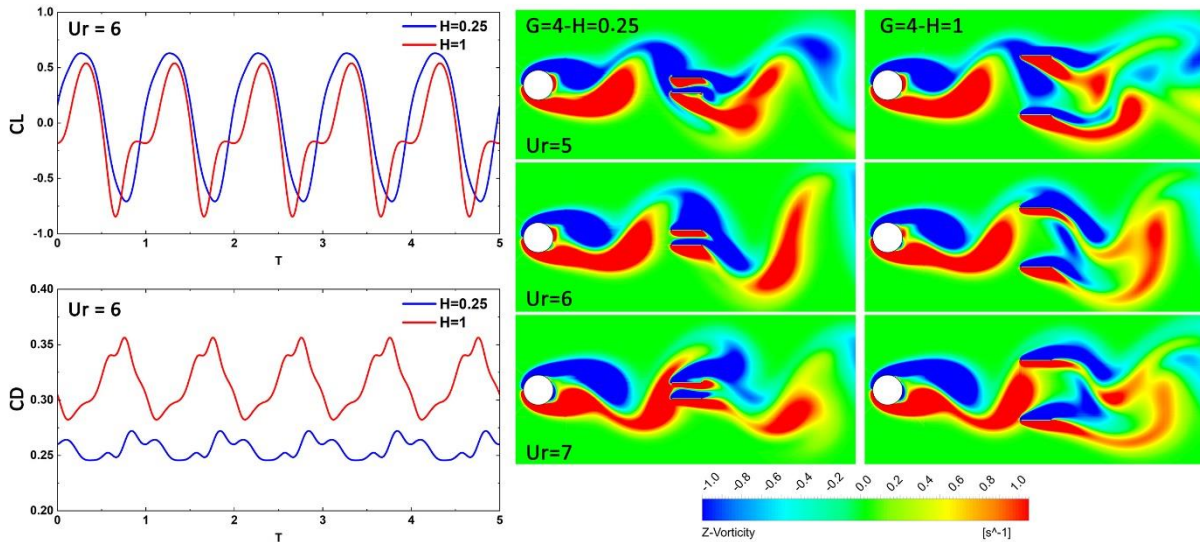
601 cylinder at all reduced velocities, without appreciably affecting the wake of the cylinder.
 602 Increased vibration amplitude at $Ur=6$ leads to decreased lift and higher mean drag coefficient
 603 of the plates due to hydrodynamic forces. The results for the second case ($G=4-H=1$) are
 604 shown in the figure. 21.



605 **FIG. 21) FIV of the parallel flat plates for $G=4-H=1$ in a range of reduced velocities**

606 In this situation, the vibration amplitude and phase difference show similar features to the
 607 previous case, with low amplitude recorded at $Ur=2-5$, a peak A^* at $Ur=6$, and a gradual
 608 decrease at higher velocities. The zenith amplitude is about 0.4, lower than the value
 609 observed for $H=0.25$. The phase difference displays a similar pattern, with a significant change
 610 occurring at $Ur=7$. The parallel plates vibrate synchronously with the vortex shedding
 611 frequency of the upstream cylinder at all tested velocities. The study of hydrodynamic forces
 612 shows that the CL_{rms} and CD_{mean} of the cylinder exhibit similar trends and magnitudes, but the
 613 plates exhibit higher mean drag values when placed at $H=1$. The lift coefficient shows small
 614 fluctuations, with higher values observed for lower reduced velocities of $Ur=2-5$.

615 Due to the similar system responses, the wake structure is compared for both configurations
 616 in the reduced velocities around the amplitude peak (FIG. 22). The variation of lift and drag
 617 coefficient for the upper plate are also presented for five cycles of vibration at $Ur=6$.

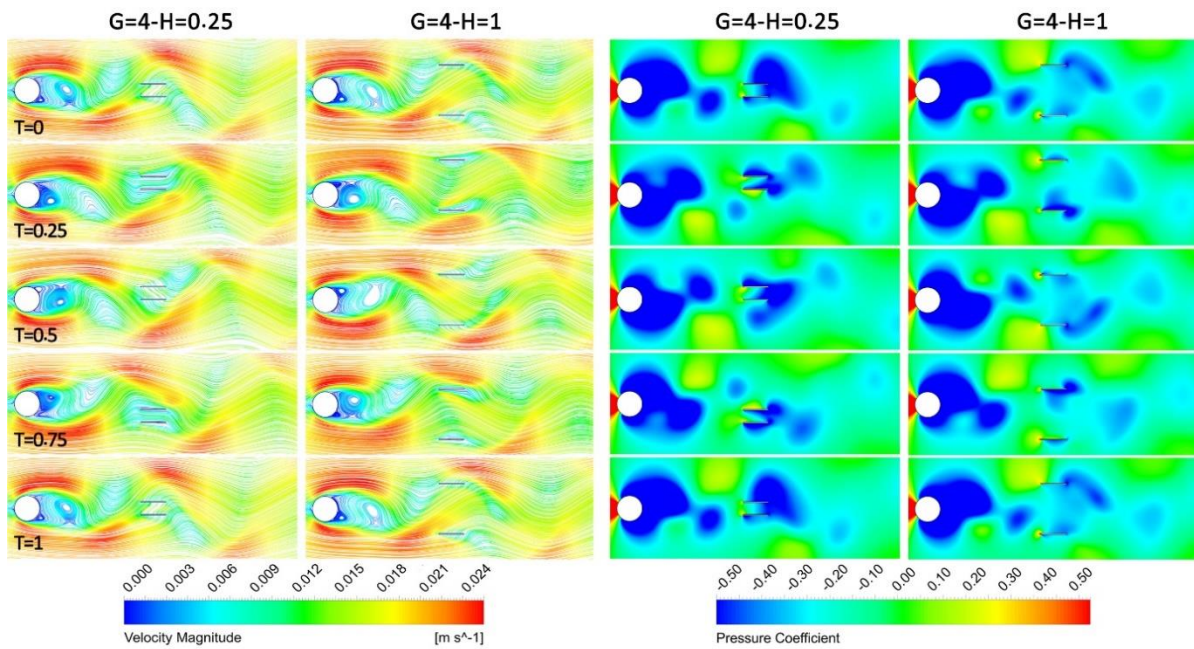


618 FIG. 22) Variation of lift and drag coefficient at $Ur=6$ and Instantaneous vorticity contours in different reduced velocities
 619 for $G=4-H=0.25, 1$

620 The wake patterns show similar features, with the plates regularly encountering the shedding
 621 vortices in each case. At $Ur=7$, a phase shift is detected as the plates interact with the vortex
 622 coming from the lower shear layer. During one vibratory cycle at $H=1$, each plate faces the
 623 vortex from its corresponding side (the top plate with the upper shear layer and vice versa),
 624 with no involvement from the other plate. When the gap is reduced to $H=0.25$, both plates
 625 contact with the oncoming vortices at the same time.

626 At $H=0.25$, the lift coefficient follows a sinusoidal pattern as shown in FIG. 22. However, at
 627 $H=1$, a secondary frequency appears, causing fluctuation in the lift coefficient. Increasing the
 628 vertical separation leads to a decrease in the average lift coefficient. The Cl_{mean} is -0.1 when
 629 the separation is increased, whereas it was 0.05 when the separation was at $H=0.25$.
 630 Analyzing CD_{mean} shows that a greater vertical distance results in higher average values and
 631 greater variations in the vibrating plates. Figure 23 illustrates the pressure coefficient and

632 streamline contours for a single vibratory cycle to better understand flow dynamics and
 633 determine the causes of these occurrences.



634 **FIG. 23) Streamlines and pressure coefficient contours for one cycle of vibration for cases $G=4-H=0.25, 1$**
 635 An analysis of the pressure coefficient contours shows that when the parallel plates are closer
 636 to the wake centerline ($H=0.25$), the vortices move past the downstream objects in a periodic
 637 manner through the upper or lower sides of both plates, supporting the sinusoidal pattern
 638 seen in the lift coefficient. By increasing the vertical distance to $H=1$, the vortices created
 639 move through the space between the parallel plates, resulting in a low-pressure area forming
 640 surrounding each plate. This occurrence leads to the appearance of a secondary frequency,
 641 causing a disturbance in the lift coefficient, as shown in FIG. 22.

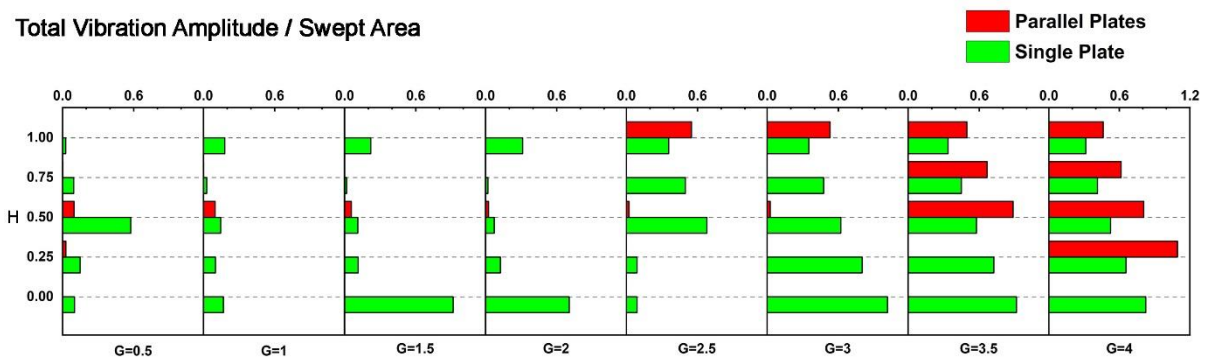
642 An in-depth analysis shows that the maximum amplitude decreases as the vertical spacing is
 643 changed to $H=1$. Two main causes are responsible for this transformation. The increased
 644 space requires the plates to interact with the external flow, which is characterized by higher
 645 energy levels. The external flow creates resistance, pushing the plates towards the centerline
 646 of the wake. The shedding vortices passing over the midway of the parallel plates apply a

647 centripetal force on the plates simultaneously. The combined forces result in a diverging
 648 mean lift coefficient, with a negative value for the upper plate and a positive value for the
 649 lower plate, leading to a reduction in vibratory amplitude. On the other hand, the opposite
 650 wake pattern caused by the vortices at $H=0.25$ results in an increased CL_{mean} , pushing the
 651 plates outward and leading to a higher vibratory intensity.

652 There is a significant rise in the drag coefficient in both setups when the Reynolds number is
 653 6. At this velocity, the plates experience increased vibration amplitude due to the flow from
 654 outer layers with higher velocity. This interaction results in heightened drag forces on the
 655 plates with a Reynolds number of 6. Increasing the vertical spacing to $H=1$, as shown in FIG.
 656 22, results in a higher mean drag coefficient and more fluctuations. The increased average
 657 drag can be explained by the plates' close closeness to the external layers, which amplifies
 658 the drag forces. The significant vibration amplitude leads the plates to oscillate back and forth
 659 in the wake zone, intensifying the fluctuations.

660 3.3 Configurations comparison

661 This study aims to investigate different setups for harvesting energy from the turbulent flow
 662 behind a cylinder. The total amplitude per swept area for each setup shown in FIG. 24 has
 663 been calculated using equation (4).

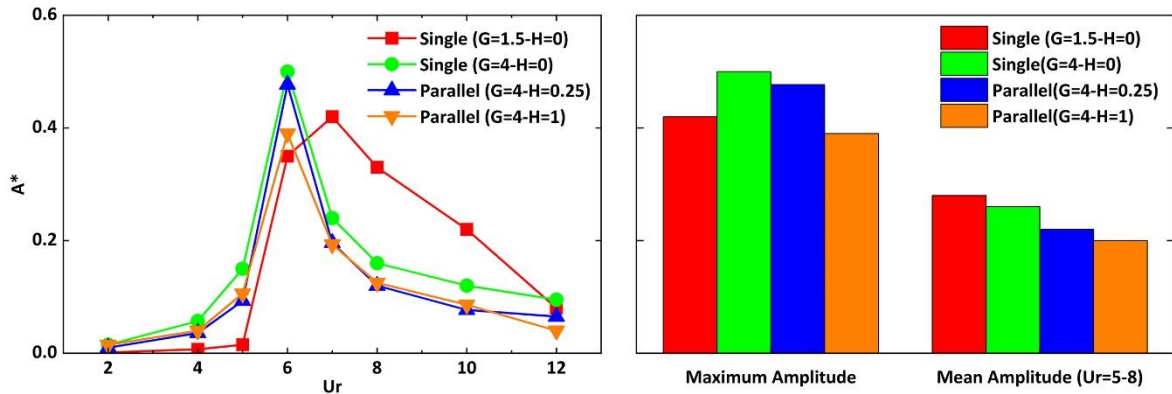


664 FIG. 24) Total vibration amplitude per swept area in different configurations at $Ur=6$

665 The chart being discussed shows significant similarities to the vibration amplitude shown in
666 FIG. 5 and 17. Similar to the vibration amplitude, a significant value is consistently visible for
667 the single plate over almost all horizontal range. At a spacing of $G=0.5-H=0.5$, the plate shows
668 a notable Vibration per Swept Area (VSA) with an amplitude of $A^*=0.18$. The narrow
669 horizontal separation and smaller swept area lead to a high VSA, despite the amplitude not
670 being significant. The increased values at $G=1.5, 2-H=0$ correspond to bigger amplitudes of
671 $A^*=0.32$ and 0.28 on the plate due to a similar reasoning. Positioning the plate near the wake
672 centerline reduces the swept area, but a bigger amplitude results in much higher VSA values.
673 A decrease in VSA is observed as the single plate's vertical alignment increases due to broader
674 horizontal separations and expanding swept regions. The scenario $G=3-H=0$ exhibits the
675 highest VSA for the single plate among all configurations, with $VSA \approx 1$.

676 In areas without vibration, the VSA is considered insignificant for parallel plates because of
677 the low amplitudes. Parallel plates can demonstrate a greater Voltage Standing Wave Ratio
678 (VSA) when the amplitude is significant, compared to similar setups with a single plate.
679 Although these situations have lower vibratory amplitudes and greater swept areas, using
680 independently vibrating entities results in a better VSA. Comparing single and parallel plates,
681 it is evident that the configuration $G=4-H=0.25$ with parallel plates obtains the highest
682 vibration amplitude per swept area among all configurations ($VSA \approx 1.1$).

683 The vibration amplitude of potential situations is thoroughly compared by assessing both
684 maximum and average amplitudes, as shown in FIG. 25.



685 FIG. 25) Comparing the mean and maximum vibration amplitudes for cases with harvesting energy potential.

686 Vibration in a range of reduced velocities for the $G=1.5-H=0$ configuration with a single plate
 687 shows unique characteristics, reaching maximum amplitude at $Ur=7$ instead of $Ur=6$. The
 688 amplitude decreases linearly, leading to higher values of A^* at lower velocities. This distinctive
 689 behavior is also evident in the highest achievable and average amplitude of the instances.
 690 While the single plate positioned at $G=4-H=0$ recorded the greatest A^* among all
 691 configurations, the average amplitude for $G=1.5-H=0$ is considered to be the most significant.
 692 Ultimately, the research shows that both setups, $G=1.5-H=0$ (single) and $G=4-H=0.25$
 693 (parallel), have significant promise for an energy harvesting system.

694 4 Conclusion

695 Numerical simulations were performed to investigate Wake-Induced Vibrations (WIV) of
 696 single or parallel flat plates, aiming to enhance understanding of vortex-plate interactions in
 697 the wake of a circular cylinder. This investigation was then widened to uncover combinations
 698 suitable to energy harvesting. Several plate configurations were examined to evaluate their
 699 impact on the upstream cylinder, expanding the analysis to cover a range of reduced velocities
 700 for specific situations. The Reynolds number, based on the free-stream velocity and the
 701 cylinder's diameter, was kept at 100. The entities had a low mass ratio of 10, and the damping
 702 effect was considered negligible ($\zeta=0$). The study culminates in the following conclusions:

703 1. The examination of WIV introduces novel perspectives as a single flat plate exhibits a
704 pronounced dependency on position within the wake of a stationary cylinder. Historically
705 seen as a passive flow control device, the flat plate begins to show significant vibrations once
706 it gets too close to a critical point. Elevating the plate's vertical position and distancing it from
707 the wake centerline alters the wake dynamics and vibratory mechanism, yielding dual
708 outcomes. A gradual reduction in the highest level of amplitude is observed, along with a
709 shrinking of the non-vibratory area, which promotes improved mixing and interaction in the
710 shear layer. The investigation also questions previous ideas about the plate's regulatory role
711 in the near-wake, showing that small horizontal separations can cause vibration. These
712 amplitudes are typically lower than those seen in wider gaps due to the unique vortex
713 formation and flow field structure in the nearby wake. Extending the analysis to a wider
714 velocity range shows that there is no lock-in regime, deviating from the behavior of circular
715 cylinders. Significant amplitudes only occur when the shedding frequency is close to the
716 natural frequency.

717 2. The two functions of parallel plates in the analyzed setups establish an important limit for
718 vibratory amplitude. Across an extensive range of separations, the plates remain inert, serving
719 as an effective flow modulation mechanism for the cylinder, manifesting two distinct
720 actions—either separating or bifurcating the shear layers, contingent on the vertical
721 separation. An abbreviated non-vibratory zone also materializes for parallel plates as the
722 vertical gap widens. The vibratory mechanism itself varies with the vertical separation,
723 encountering periodic shedding vortices at narrower gaps and confronting vortices
724 emanating from the shear layer on the same side at wider gaps. Lock-in is not present in the
725 chosen scenarios at different reduced velocities, but significant vibration amplitudes occur
726 when the frequencies of vortex shedding closely match the natural frequency.

727 3. Comparing single plates to parallel plates in terms of vibratory amplitude and spatial
728 occupation provides new perspectives for designing energy harvesting devices. Specifically,
729 two distinct configurations are identified, each offering unique benefits. The first plate,
730 positioned at $G=1.5-H=0$, exhibits moderate vibratory characteristics with a decreased
731 horizontal distance and a smaller area of occupancy. On the other hand, the second
732 configuration, with parallel plates placed at a distance downstream at $G=4-H=0.25$,
733 demonstrates the benefits of using separate vibrating components.

734 The present study may be extended further to explore the effect of different cross-sections
735 for the upstream cylinder to achieve higher vibration amplitudes. Additionally, future
736 research could investigate increasing the Re number to the subcritical regime and examining
737 the impact of changing variables such as the mass and damping ratios to examine the results
738 under more realistic conditions.

739 Declaration of competing interest

740 The authors declare that they have no known competing financial interests or personal
741 relationships that could have appeared to influence the work reported in this paper.

742 Acknowledgment

743 "The authors acknowledge the use of the UCL Grace High Performance Computing Facility
744 (Kathleen@UCL), and associated support services, in the completion of this work.

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